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# Wavelet transforms and their applications for ITER

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*ITER International School  
Aix-en-Provence  
26° August 2014*





晴海亭

越平



# Outline

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## 1. *Wavelet transforms*

*Integral transforms*

*Continuous wavelet transform*

*Fast orthogonal wavelet transform*

## 2. *Applications*

*Extraction of coherent structures  
out of turbulent flows*

## 3. *Proposals for ITER*

*Analyzing signals measured in the edge plasma*

*Denoising fast camera visible light movies*

*Improving PIC codes*

# Integral transforms

Analysis

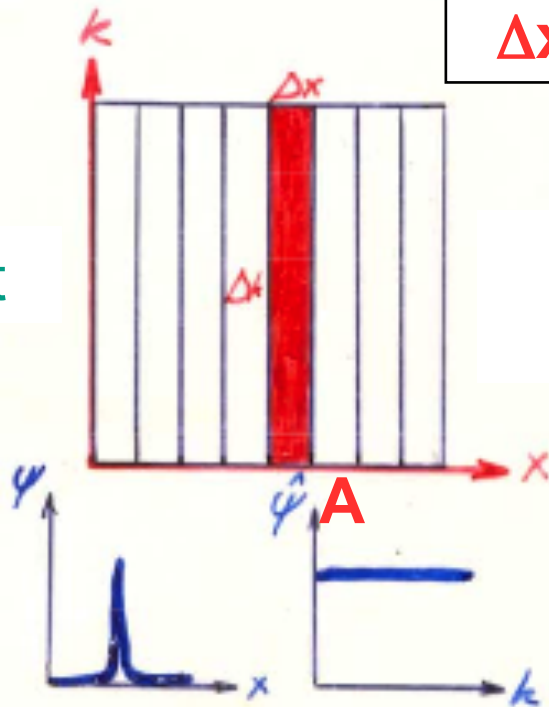
$$T_f(k) = \int f(x) \psi_k(x) dx$$

Synthesis

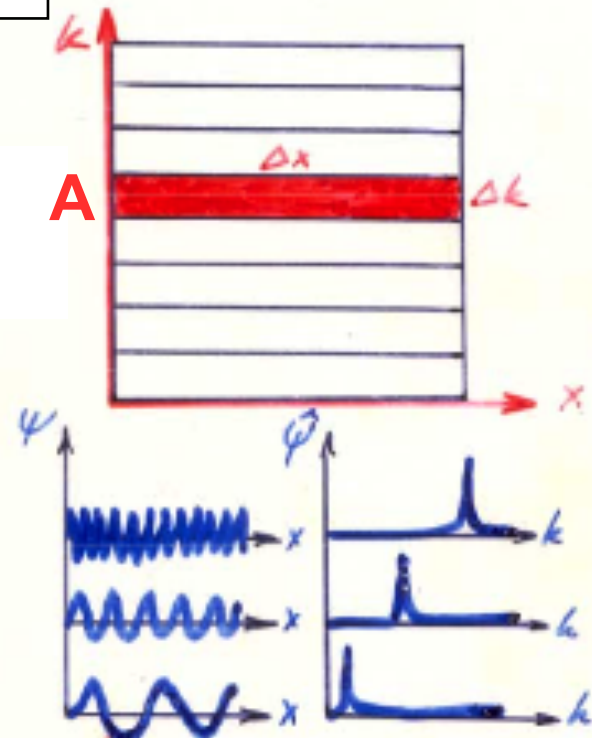
$$f(x) = \frac{1}{c} \int T_f(k) \psi_x(k) d\mu(k)$$

$$\Delta x \Delta k = A$$

Gridpoint



Fourier  
(1807)





# Continuous Fourier transform

---

$$f(x) \in L^1(\mathbb{R}) \cap L^2(\mathbb{R})$$

*Analysis*

$$\hat{f}(k) = \int_{-\infty}^{\infty} f(x) e^{-i2\pi k \cdot x} dx$$

*Synthesis*

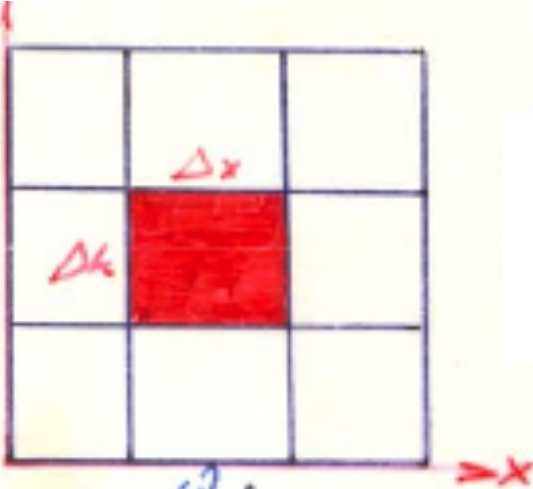
$$f(x) = \int_{-\infty}^{\infty} \hat{f}(k) e^{i2\pi k \cdot x} dk$$

*Parseval's identity*

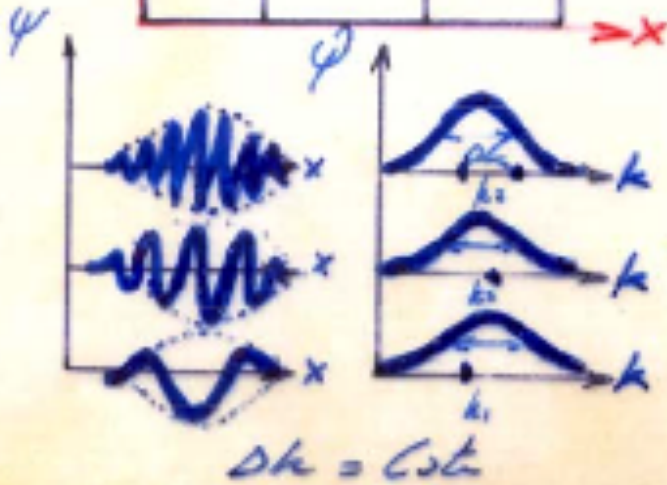
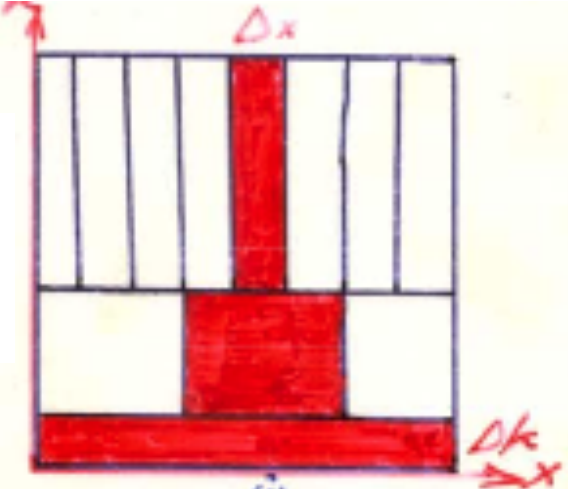
$$\int_{-\infty}^{\infty} f_1(x) \cdot f_2(x) dx = \int_{-\infty}^{\infty} \hat{f}_1(k) \cdot \hat{f}_2(-k) dk$$

# Optimal phase space tiling

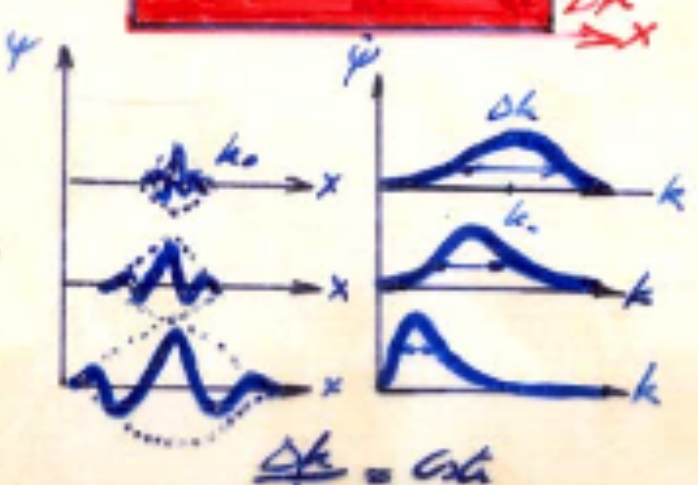
Gabor  
(1946)



Wavelets  
(1984)



Balian's  
destruction  
(1981)



Space-wavenumber  
representation

$\Delta x \Delta k =$   
information atom

Space-scale  
representation



# Choice of the analyzing wavelet

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*Admissibility condition*

$$C_\psi = \int_0^\infty \frac{|\hat{\psi}(k)|^2 dk}{|k|} < \infty$$
$$\int_{-\infty}^\infty \psi(x) dx = 0 \quad \text{or} \quad \hat{\psi}(0) = 0$$

*Jean Morlet*



*Alex Grossmann*



*Analyzing wavelet family  
generated by translation ( $b$ )  
and dilation ( $a$ )*

$$\psi_{a,b}(x) = \frac{1}{\sqrt{a}} \psi \left( \frac{x-b}{a} \right)$$

Grossmann and Morlet,  
*Decomposition of Hardy functions into square  
integrable wavelets of constant shape,*  
*SIAM J. math. Anal.*, **15**(4), 723-736, 1984

# Continuous wavelet transform (CWT)

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*Analysis*

$$\tilde{f}(a, b) = \int_{-\infty}^{\infty} f(x) \psi_{a,b}^*(x) dx$$

*Synthesis*

$$f(x) = \frac{1}{C_\psi} \int_0^{\infty} \int_{-\infty}^{\infty} \tilde{f}(a, b) \psi_{a,b}(x) \frac{da db}{a^2}$$

*Parseval's identity*

$$\langle f_1, f_2 \rangle = \int_{-\infty}^{\infty} f_1(x) f_2^*(x) dx = \frac{1}{C_\psi} \int_0^{\infty} \int_{-\infty}^{\infty} \tilde{f}_1(a, b) \tilde{f}_2^*(a, b) \frac{dadb}{a^2}$$

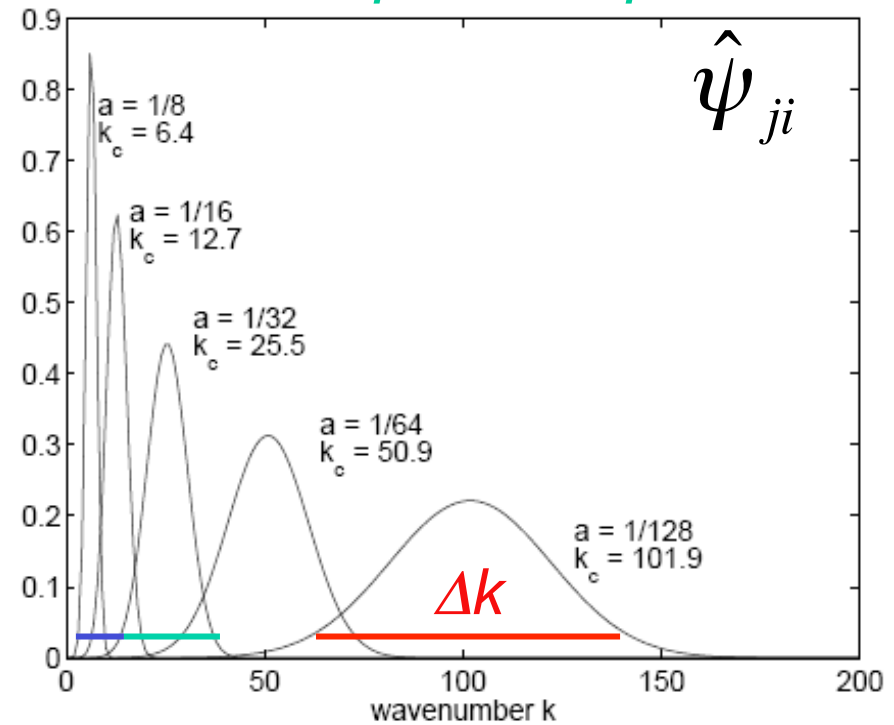
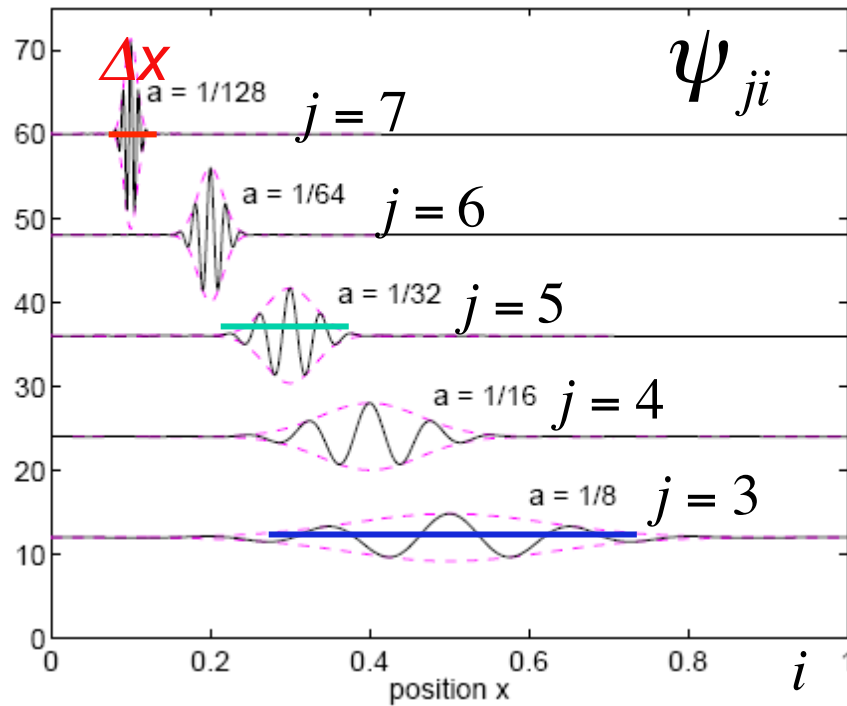


# Wavelet representation

Physical space

$$\Delta x \Delta k > C$$

Spectral space

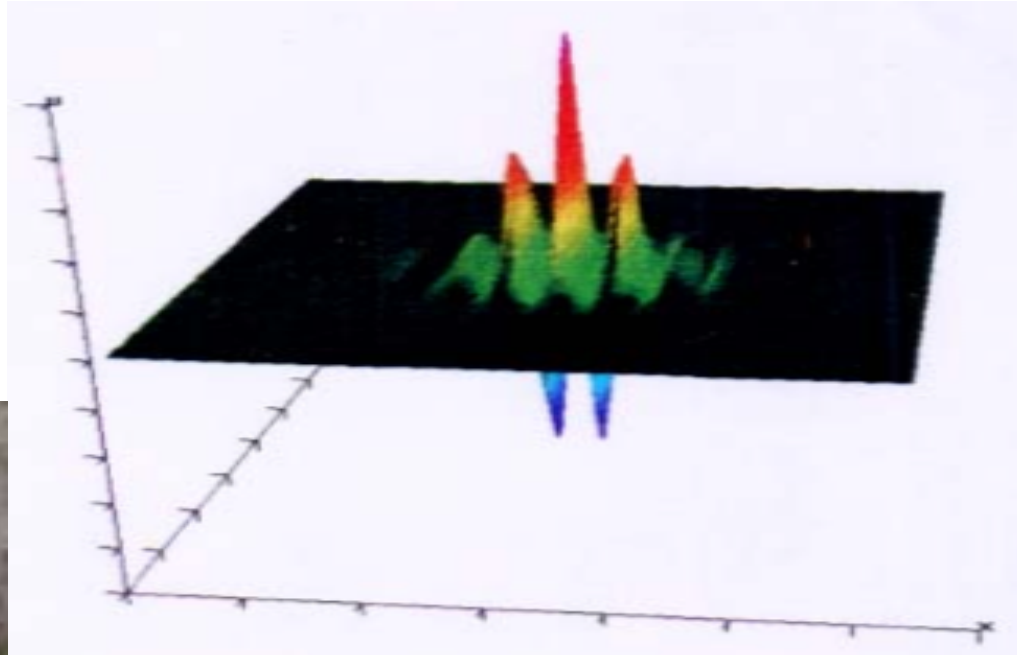


Farge,  
Wavelet transforms and  
their applications to turbulence  
*Ann. Rev. Fluid Mech.*, **92**, 1992

Farge and Schneider,  
Wavelets: application to turbulence,  
*Encyclopedia of Mathematical Physics*,  
Springer, 408-42, 2006

# 2D continuous wavelet transform

Jean-Pierre  
Antoine



Romain  
Murenzi



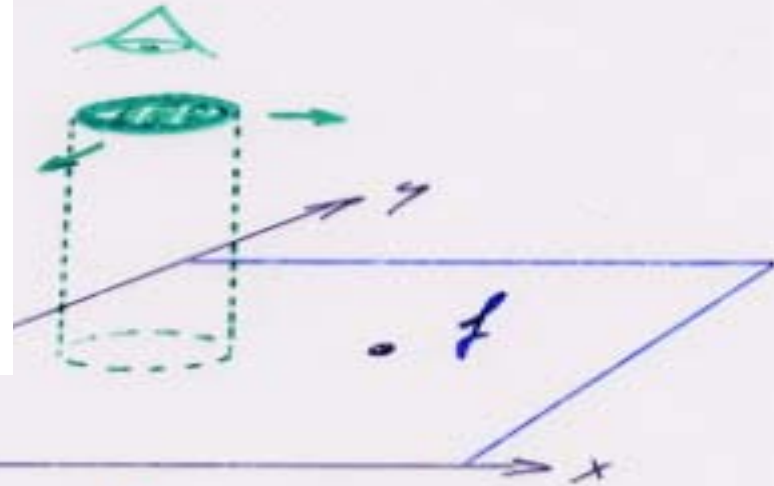
*2D Morlet mother wavelet*

*The wavelet family is generated  
by translating, dilating and rotating  
the 2D mother wavelet*



Analyzing wavelet

Field to analyze



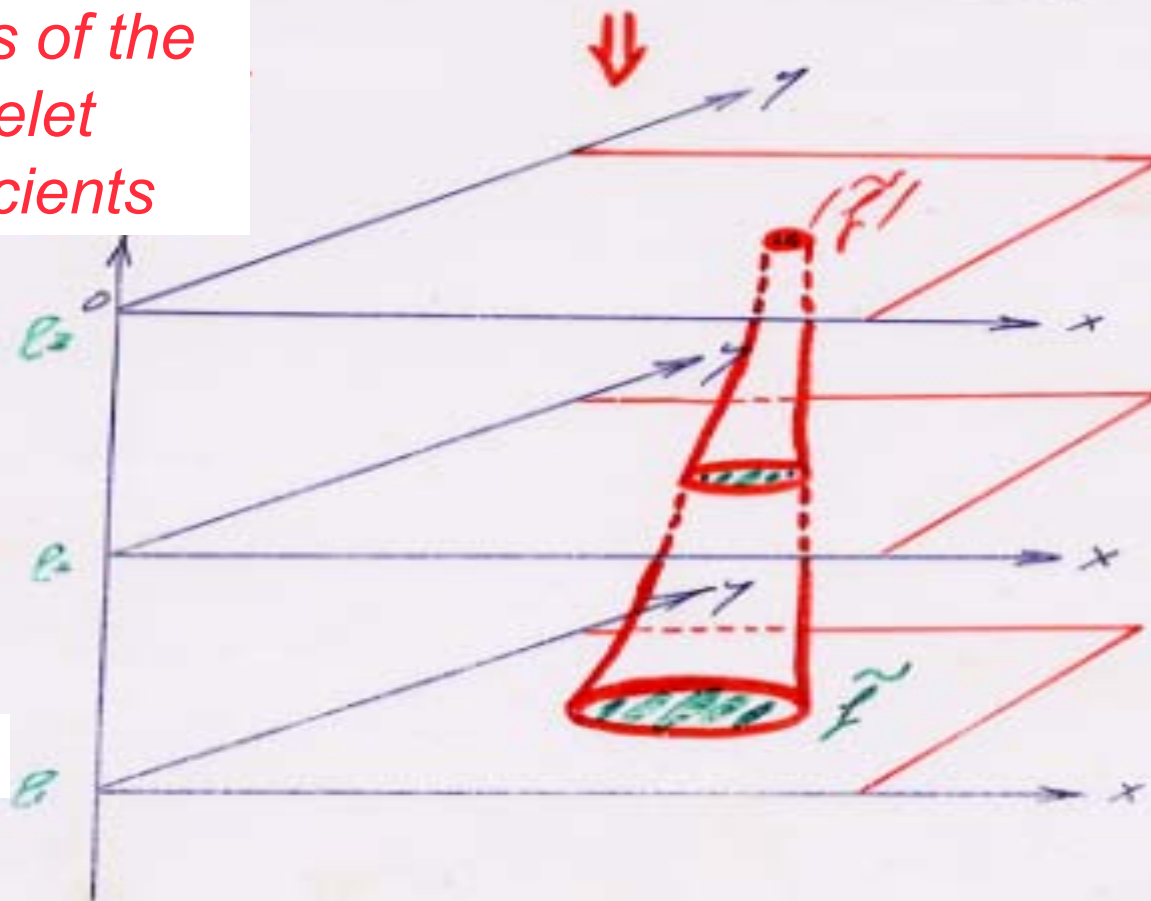
Modulus of the  
wavelet  
coefficients



Small scale



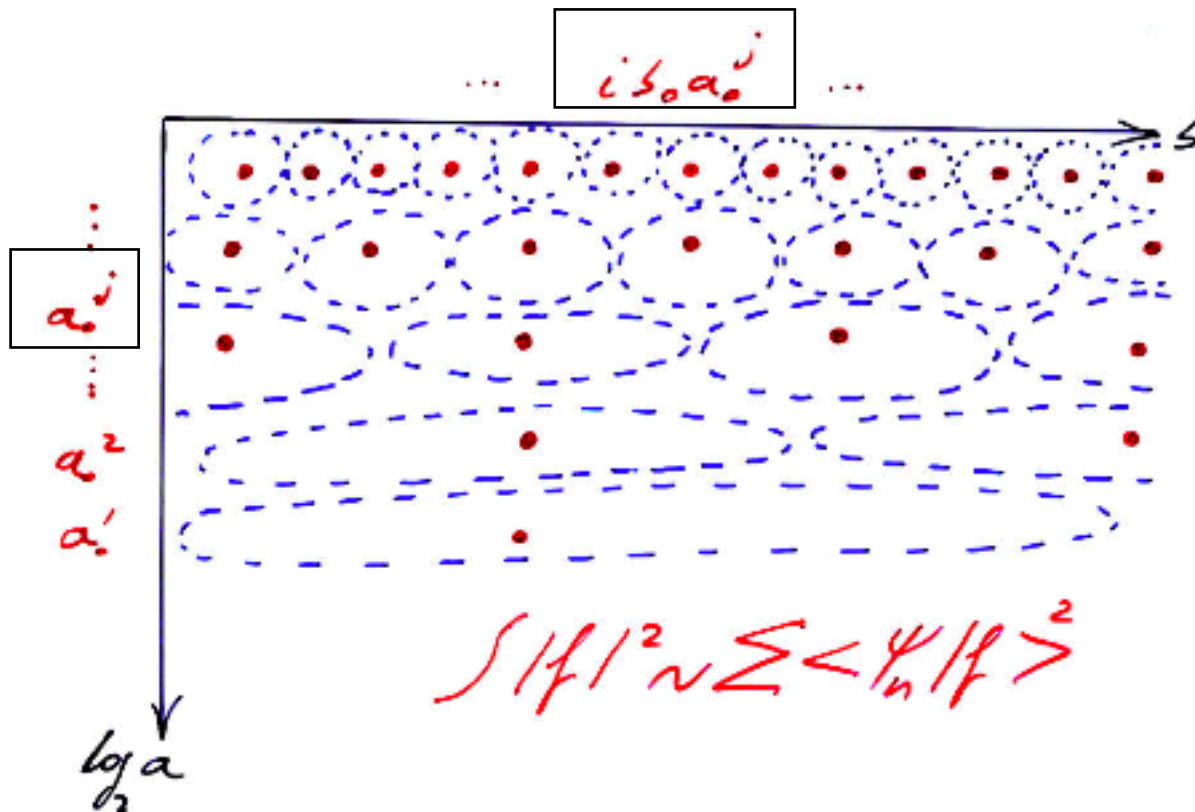
Large scale



Logarithm  
of the scale

# Wavelet frame

We can then select a finite number of wavelets restricted to a discrete grid optimally chosen such that the wavelet family associated to this grid constitutes a quasi-orthogonal basis  $\Rightarrow$  a wavelet frame



For example  
for Marr wavelet

we need

$$a_0 = 2^{1/2}$$

$$b_0 = 1/2$$

# Orthogonal wavelet transform

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Wavelet analysis :

$$\tilde{f}_{ji} = \langle \psi_{ji} | f \rangle \quad \text{with} \quad \psi_{ji} = 2^{j/2} \psi(2^j x - i)$$

Wavelet synthesis :

$$f = \sum_{ji} \langle \psi_{ji} | f \rangle \psi_{ji}$$

*A signal sampled on  $N$  points is wavelet analyzed and synthesized in  $CN$  operations if one uses compactly-supported wavelets computed from a quadratic mirror filter of length  $M$ .*

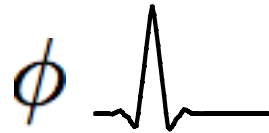


# 2D orthogonal wavelets

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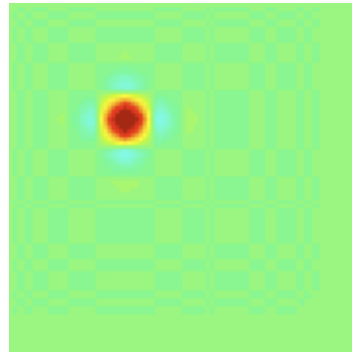
*Scaling function*

*Wavelet*



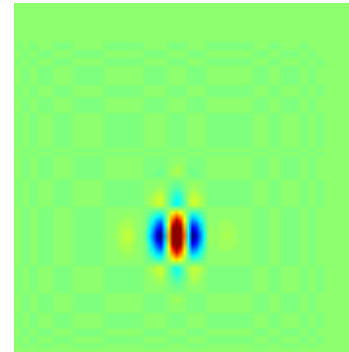
$\phi(x)\phi(y)$

Coarse  
approximation



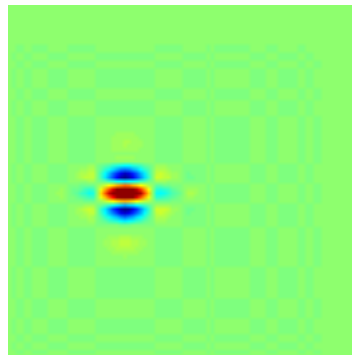
$\psi(x)\phi(y)$

Horizontal  
details



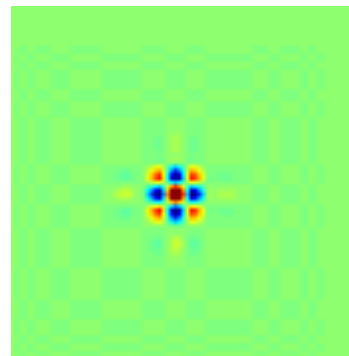
$\phi(x)\psi(y)$

Vertical  
details



$\psi(x)\psi(y)$

Diagonal  
details



# 3D orthogonal wavelets

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A 3D vector field  $v(x)$  sampled on  $N = 2^{3J}$  equidistant grid points

$\psi_\lambda(x)$  3D wavelet  $\rightarrow$  orthogonal wavelet series

$$v(x) = \sum_{\lambda \in \Lambda} \tilde{v}_\lambda \psi_\lambda(x), \quad \tilde{v}_\lambda = \langle v, \psi_\lambda \rangle$$

$$\Lambda = \{ \lambda = (j, i_n, \mu, \psi_\lambda(x)), \dots, J-1, i_n = 0, \dots, 2^j - 1, n = 1, 2, 3, \text{ and } \mu = 1, \dots, 7 \}$$

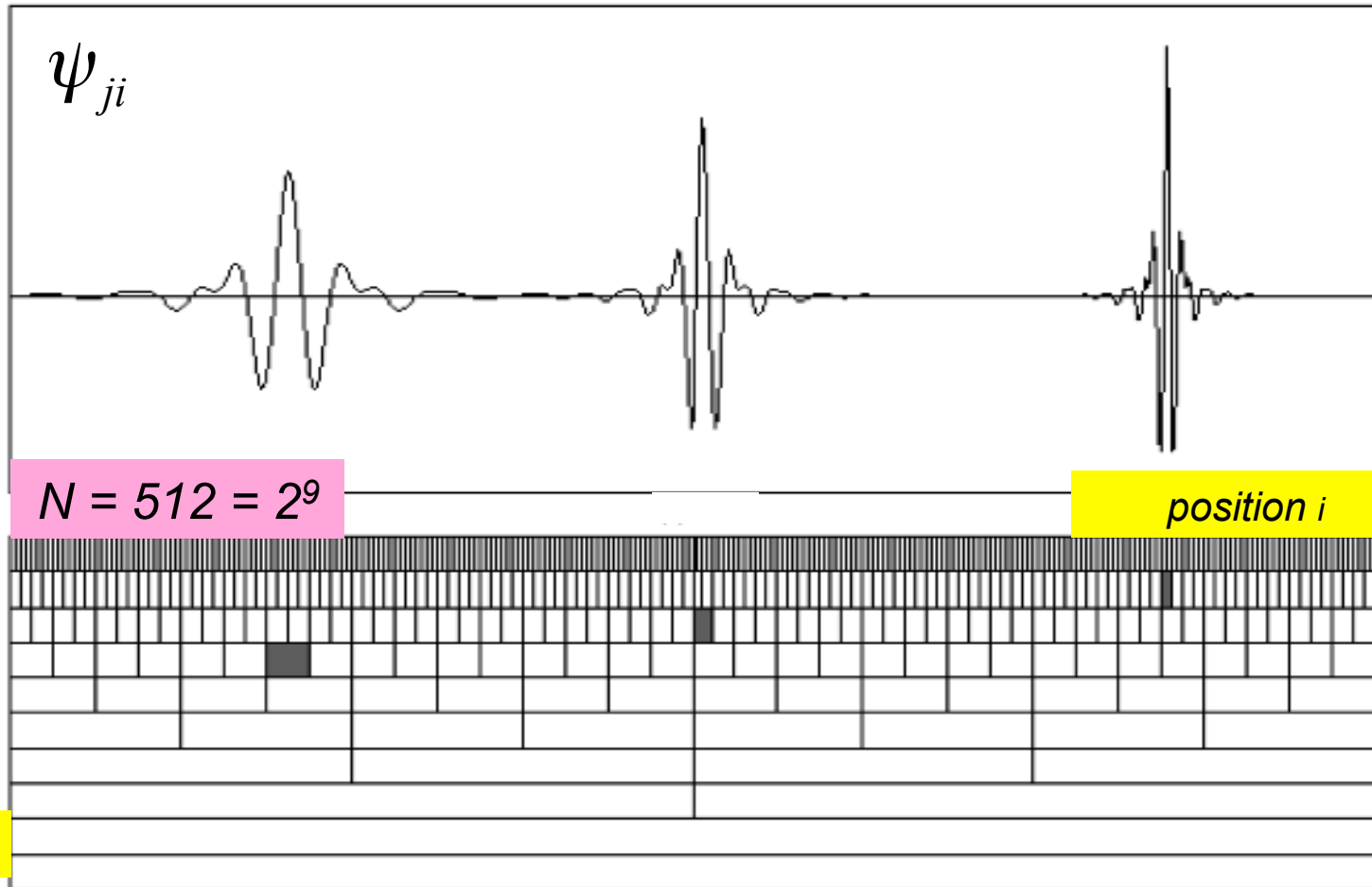
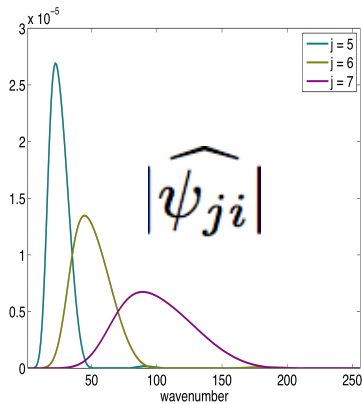
$N_j = 7 \times 2^{3j}$ , wavelet coefficients at a scale indexed by  $j$

- fast algorithm with linear complexity
- no redundancy between the coefficients

We use Coifman 12 wavelet  
compactly supported with four vanishing moments.

# Orthogonal wavelet representation

## Wavelets



## Wavelet coefficients

$$\tilde{f}_{ji} = \langle \psi_{ji} | f \rangle$$

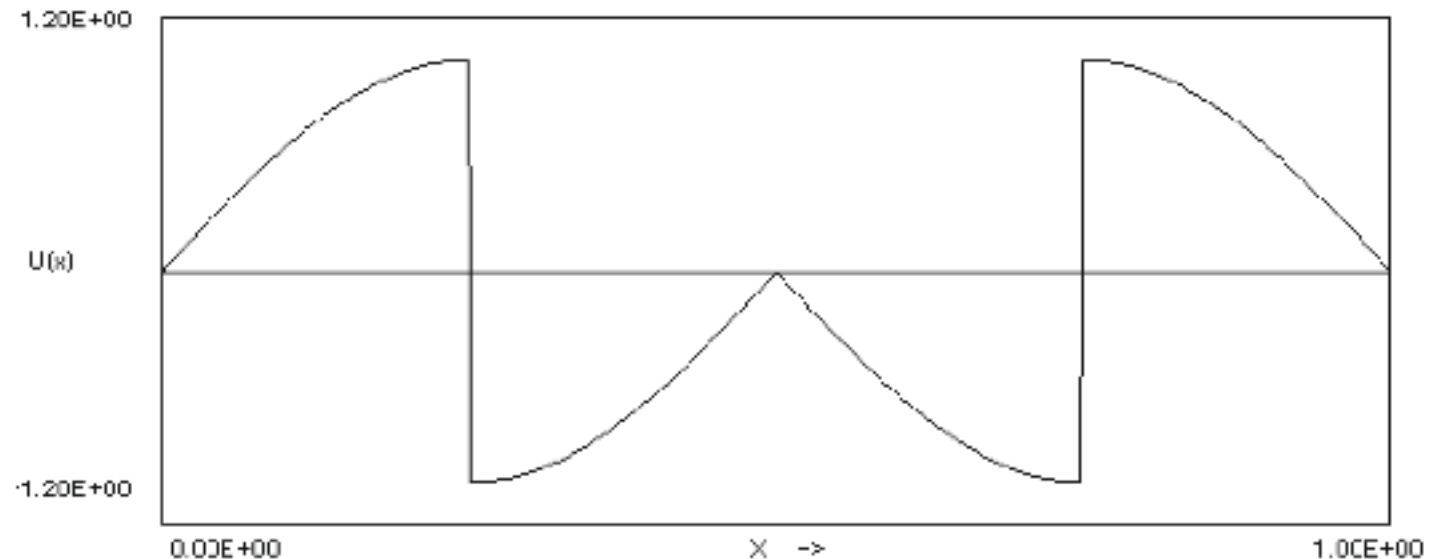
scale  $j$

Mallat,  
A wavelet tour of  
signal processing, 3rd edition,  
Academic Press, 2008

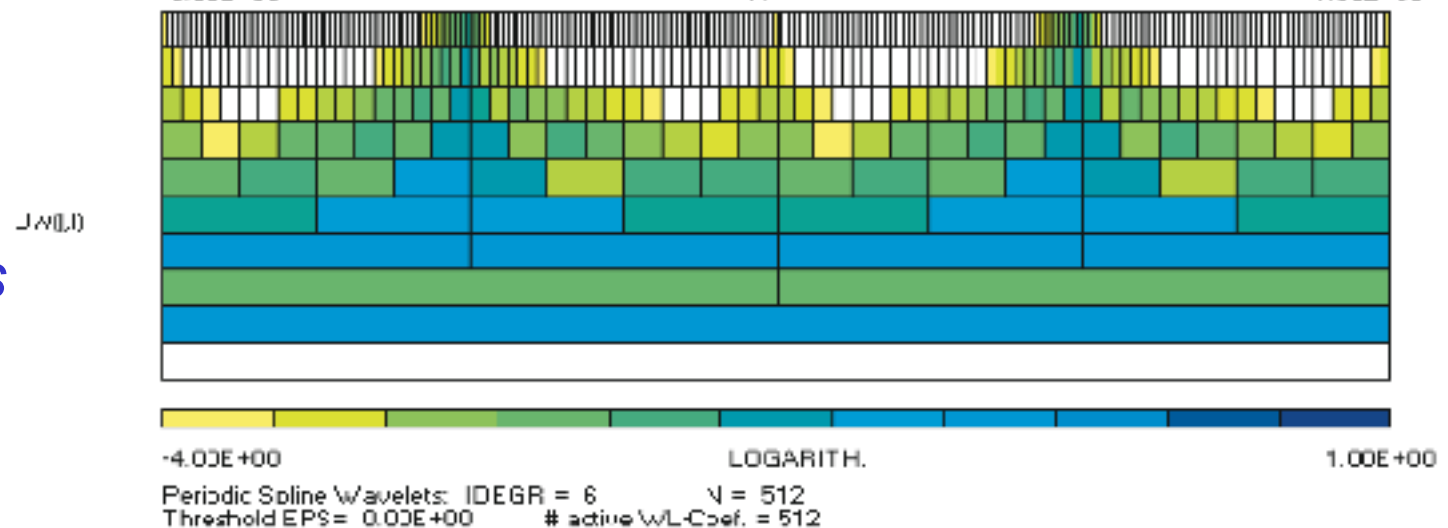


# Academic example

*Function  
sampled on  
 $N = 512$   
grid-points*



*$N = 512$   
wavelet  
coefficients*

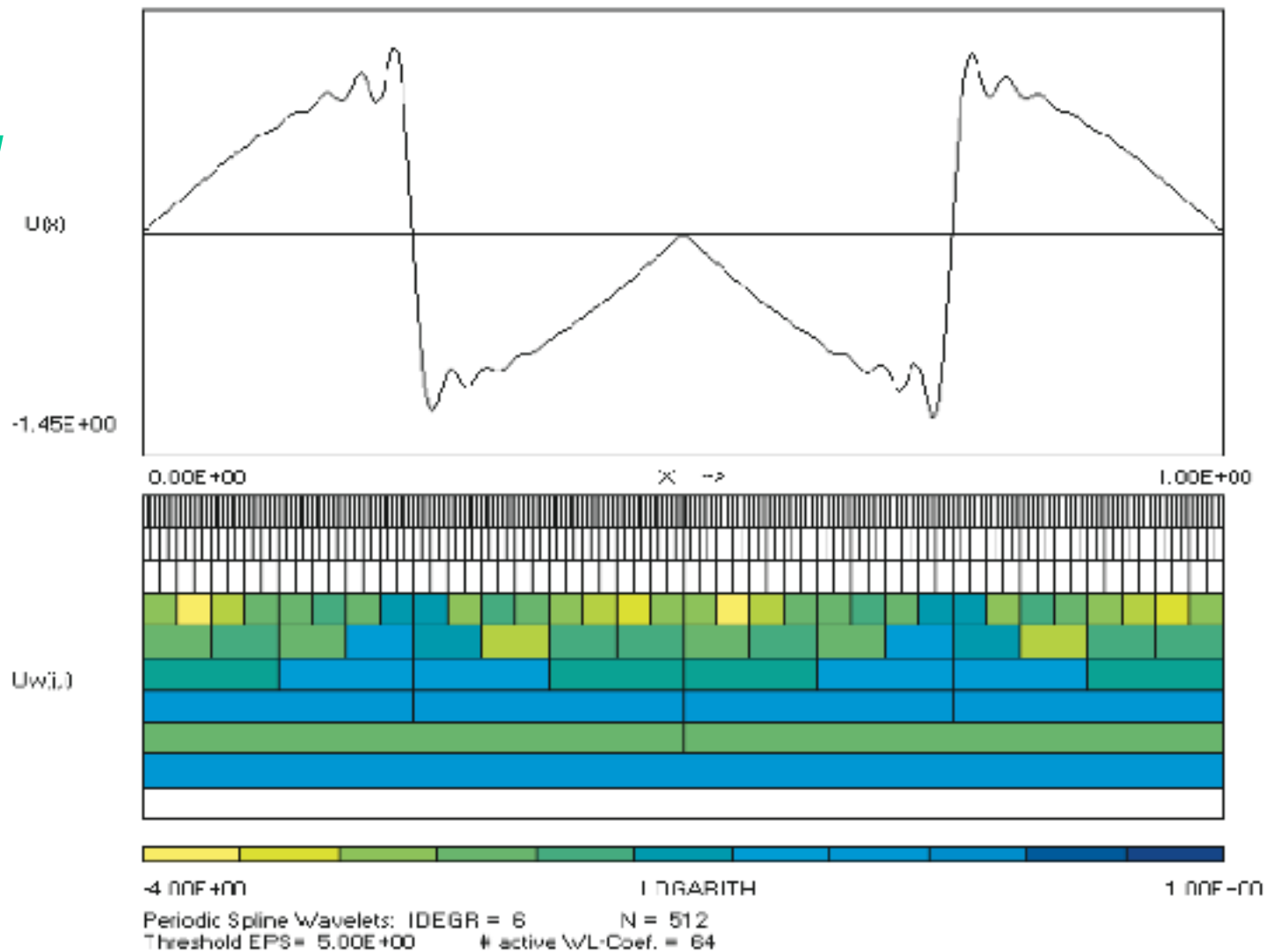


# Linear approximation

Function reconstructed from 64 wavelet coefficients

64 wavelet coefficients such that

$$j \leq 6$$

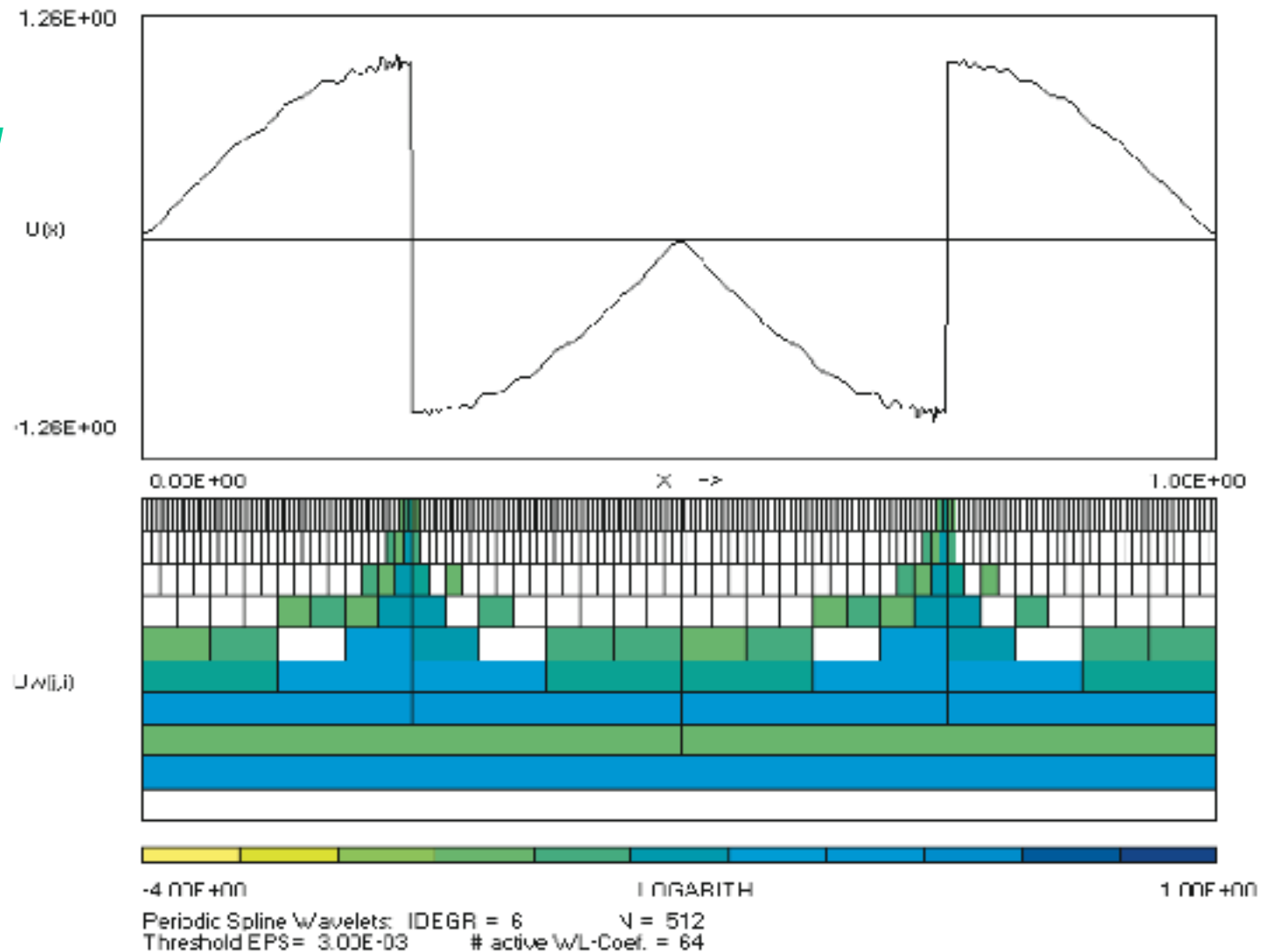


# Nonlinear approximation

Function reconstructed from 64 wavelet coefficients

64 wavelet coefficients such that

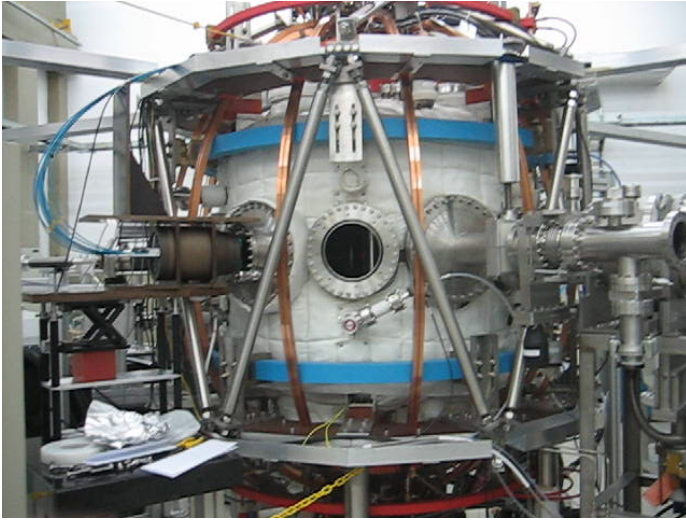
$$|\tilde{f}_{ji}| > \epsilon$$



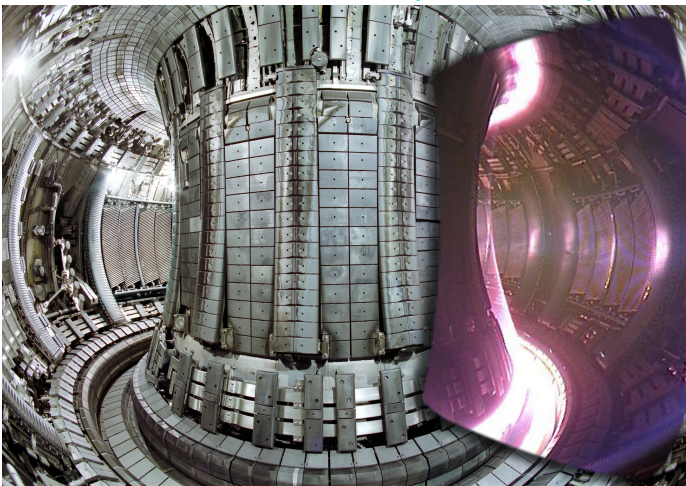


# Application to tokamaks

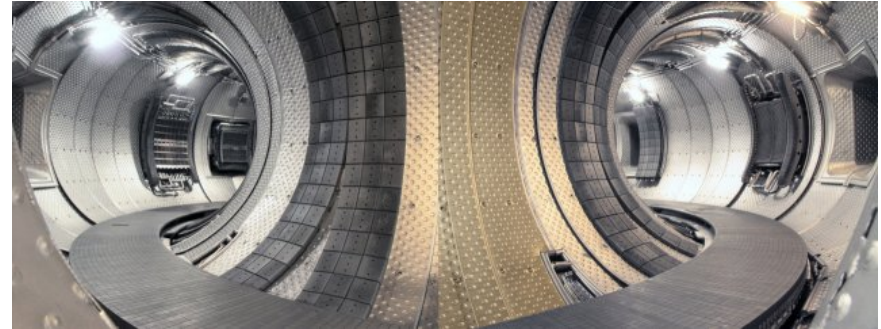
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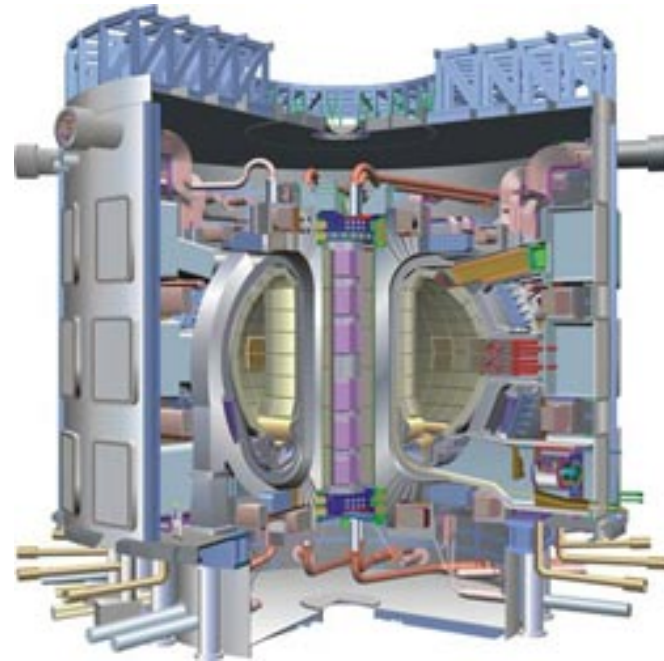
ETE, INPE (Brazil)



JET, Culham (Europe)



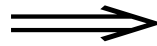
Tore-Supra, Cadarache (France)



ITER (World)

# Turbulent edge plasma in the SOL

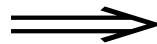
Turbulent edge plasma  
in the SOL (Scrape Off Layer),  
where there are very large density  
and temperature gradients



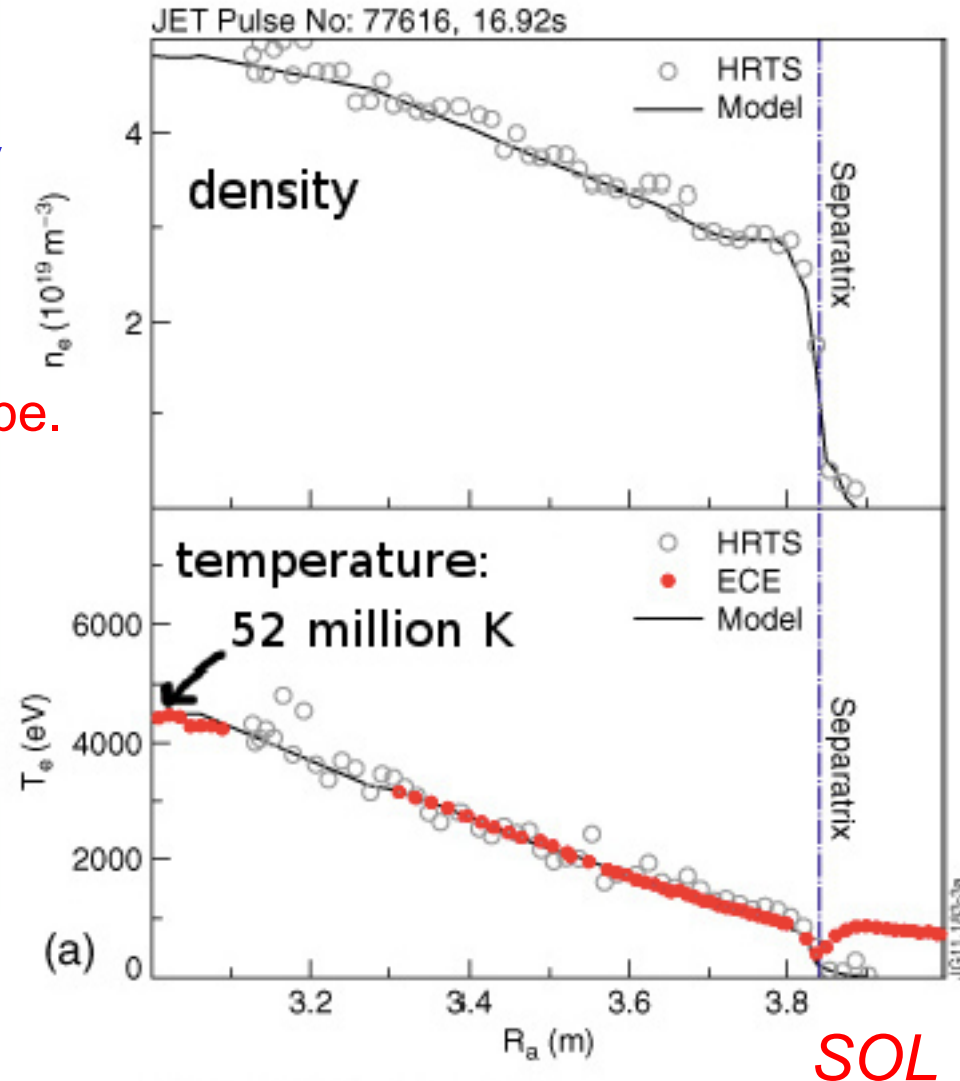
Ion density fluctuations measured  
by a fast reciprocating Langmuir probe.

Edge plasma is colder  
than core plasma, then ions  
and electrons can recombine.

Recombination and  
later desexcitation induce  
visible light emission  
(e.g., H $\alpha$  line with  $\lambda=656$  nm).



Video acquisition using  
a fast camera (40 KHz).

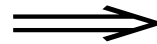


# How to extract coherent structures?

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Since there is **not yet a universal definition of coherent structures** which emerge out of turbulent fluctuations due to the nonlinear interactions, **we adopt an apophetic method** :  
**instead of defining what they are, we define what they are not.**

*For this we propose the minimal statement :*  
***'Coherent structures are not noise'***



Extracting coherent structures becomes a **denoising problem**,  
**not requiring any hypotheses on the structures** themselves  
**but only on the noise** to be eliminated.

Choosing the **simplest hypothesis** as a first guess,  
we suppose we want to eliminate an **additive Gaussian white noise**,  
and for this we use a **nonlinear wavelet filtering**.

*Farge, Schneider et al.,  
Phys. Fluids, 15 (10), 2003*

*Azzalini, Farge, Schneider,  
ACHA, 18 (2), 2005*



# Denoising using wavelets

---

Gaussian **white noise** is by definition **equidistributed** among all **modes** and the amplitude of the coefficients is given by its r.m.s., whatever the functional basis one considers.

Therefore the **coefficients of a noisy signal whose amplitudes are larger than the r.m.s. of the noise belong to the denoised signal**. This procedure corresponds to **nonlinear filtering**.

The advantage of performing such a nonlinear filtering using the wavelet representation is that **the wavelet coefficients preserve the space locality**, since wavelets are functions localized in both physical and spectral space.

Since we do not know *a priori* the r.m.s. of the noise, we have proposed an **iterative procedure** which takes as first guess the r.m.s. of the noisy signal.

Azzalini, M. F., Schneider, 2005  
*Appl. Comput. Harmonic Analysis*, **18** (2)

# Wavelet denoising algorithm

## Aphotic method :

- no hypothesis on the structures,
- *only hypothesis on the noise,*
- *simplest hypothesis as our first choice.*

## Hypothesis on the noise :

$$f_n = f_d + n$$

- $n$  Gaussian white noise,  
 $\langle f_n^2 \rangle$  variance of the noisy signal,  
 $N$  number of coefficients of  $f_n$ .

## Wavelet decomposition :

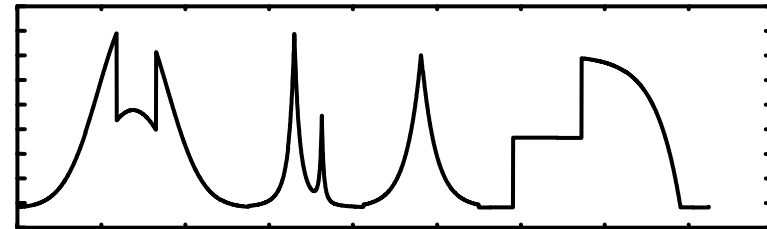
$$\tilde{f}_{ji} = \langle f | \psi_{ji} \rangle \quad \begin{array}{l} j \text{ scale,} \\ i \text{ position} \end{array}$$

## Estimation of the threshold :

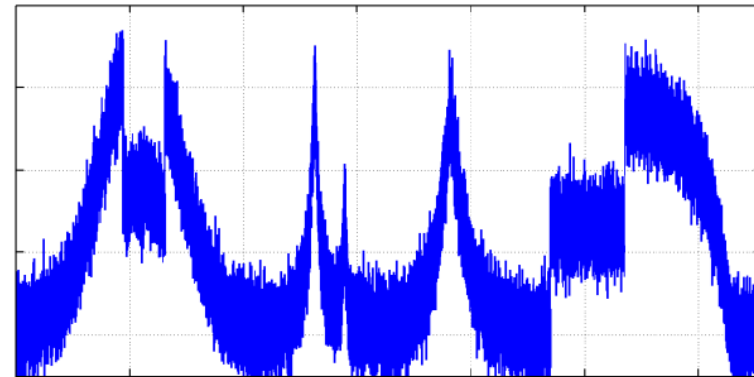
$$\varepsilon_n = \sqrt{2 \langle f_n^2 \rangle \ln(N)}$$

## Wavelet reconstruction :

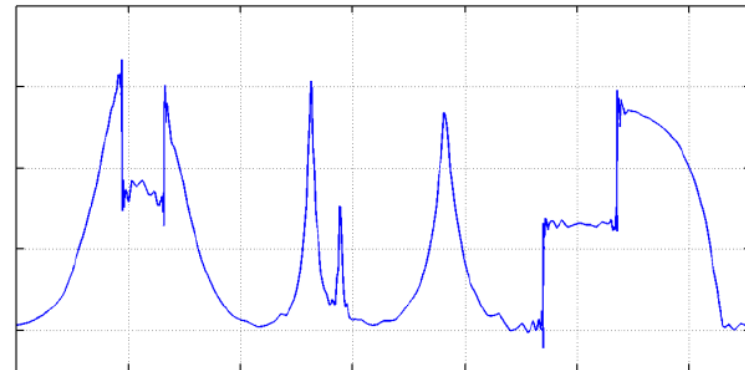
$$f_d = \sum_{j,i: |\tilde{f}_{ji}| > \varepsilon_n} \tilde{f}_{ji} \psi_{ji}$$



$f$



$f_n$



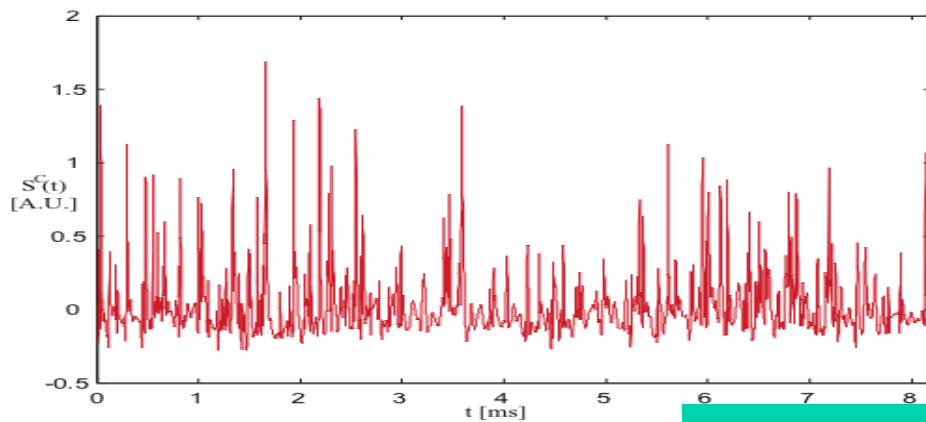
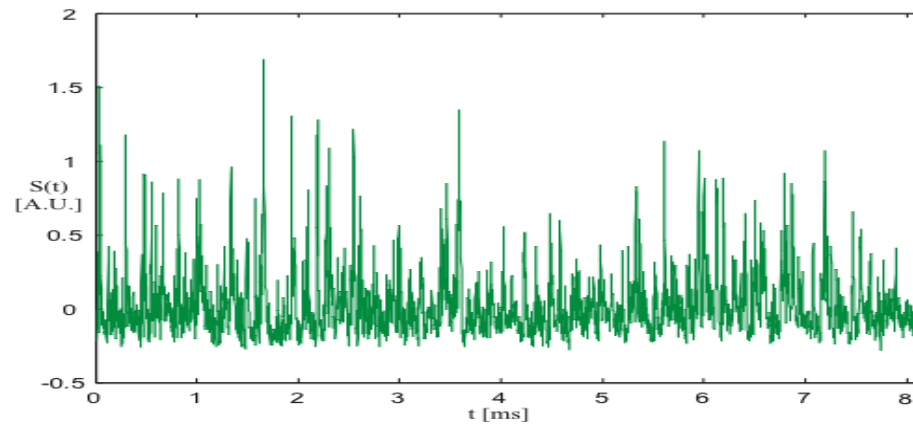
$f_d$

Donoho, Johnstone,  
Biometrika, **81**, 1994

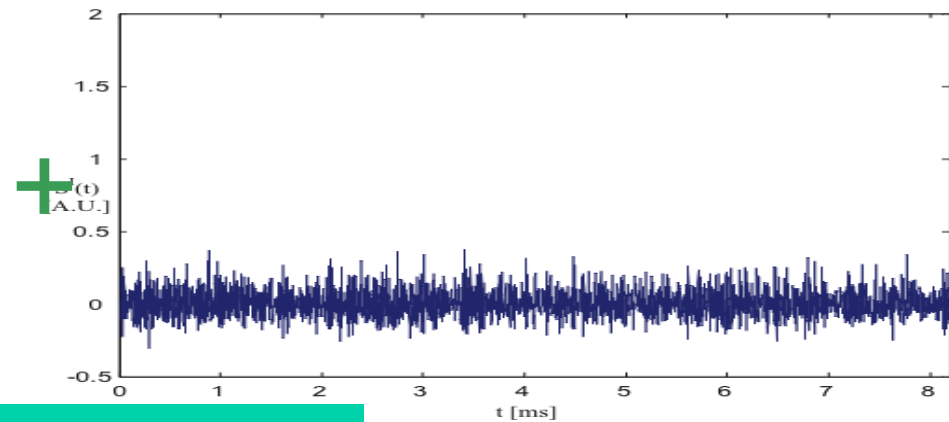
Azzalini, M. F., Schneider,  
ACHA, **18** (2), 2005

# Extraction of coherent structures SOL

Ion density fluctuations measured by a fast reciprocating Langmuir probe in the SOL of the tokamak Tore Supra  
(Pascal Devynck, Tore-Supra, CEA-Cadarache)



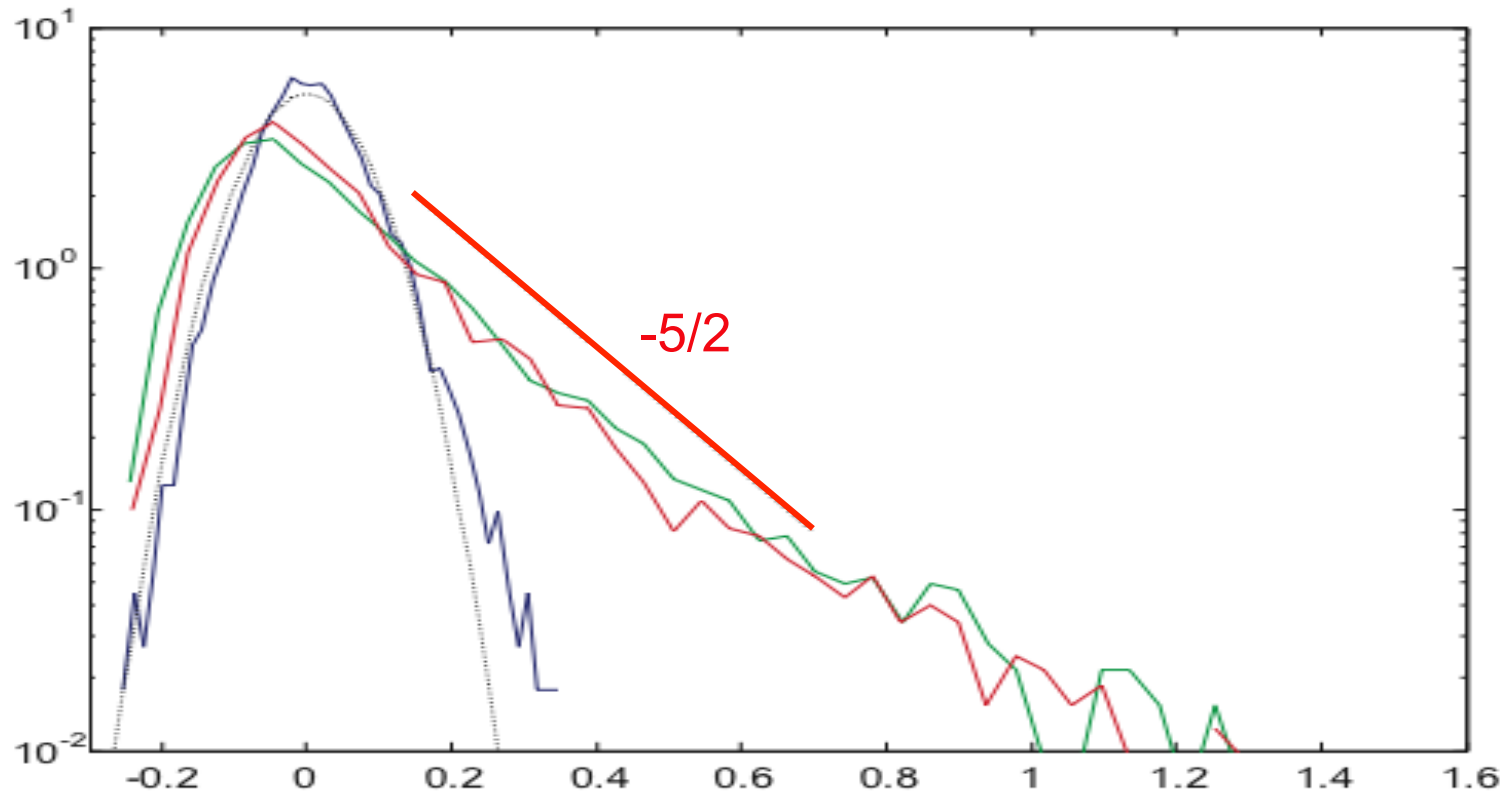
Coherent



Incoherent

Farge, Schneider & Devynck  
*Phys. Plasmas*, 13, 042304, 2006

# PDF of the density fluctuations

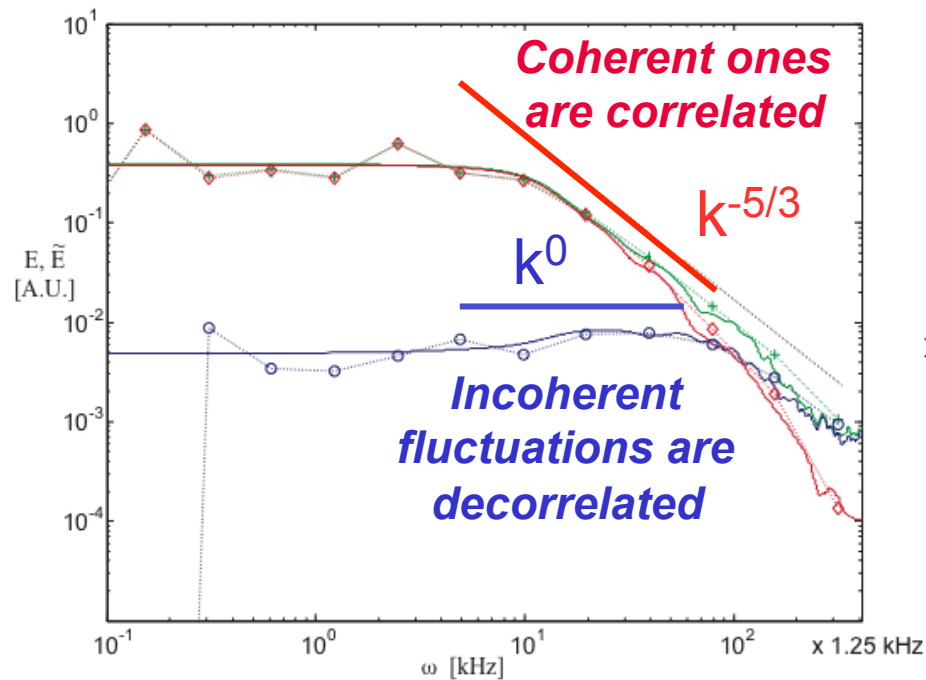


Total fluctuations = coherent + incoherent fluctuations

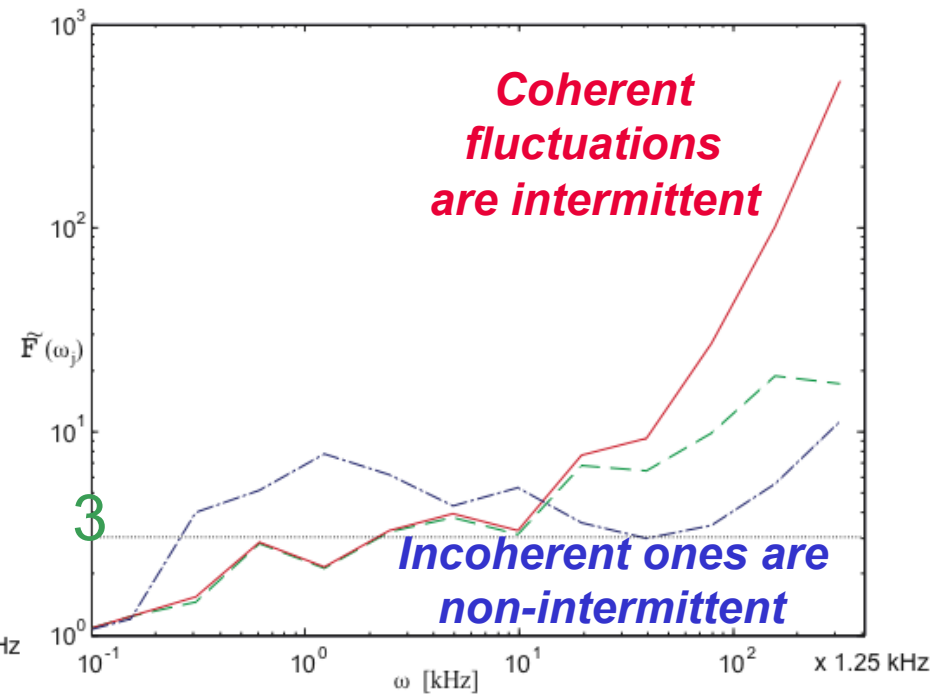
Farge, Schneider & Devynck,  
*Phys. Plasmas*, **13**, 2006

# Correlation and intermittency

**Scalogram**  
(stabilized periodogram)



**Flatness versus scale**  
(from wavelet coefficients)



Total fluctuations = coherent + incoherent fluctuations



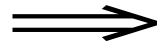
# Fast visible light camera

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A fast camera from the Nancy team (*G. Bonhomme and F. Brochard*) was installed on Tore-Supra (*N. Fedorczak and P. Monier-Garbet*).

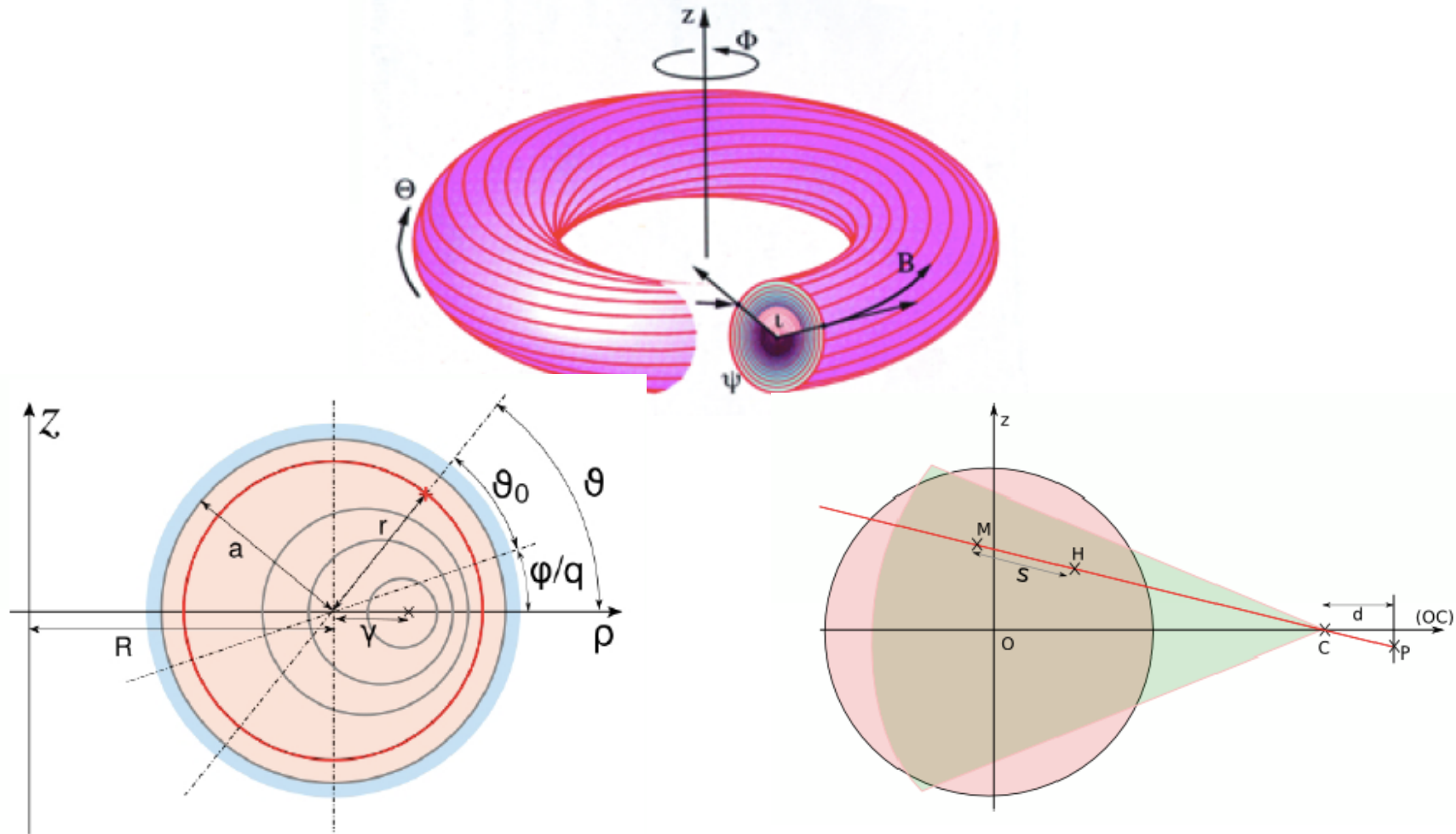
An helical Abel transform relates the plasma light emissivity  $S$  to the integral of the volume emissivity received by the camera  $I=KS$ , where  $K$  is a compact continuous operator.

Reconstruction of  $S$  from  $I$  is an inverse problem which becomes very difficult when  $S$  is corrupted by noise, then solving  $K^{-1}$  is an ill-posed problem.



Tomographic inversion using wavelet-vaguelette decomposition as an alternative to SVD (Singular Value Decomposition).

# Image tomography

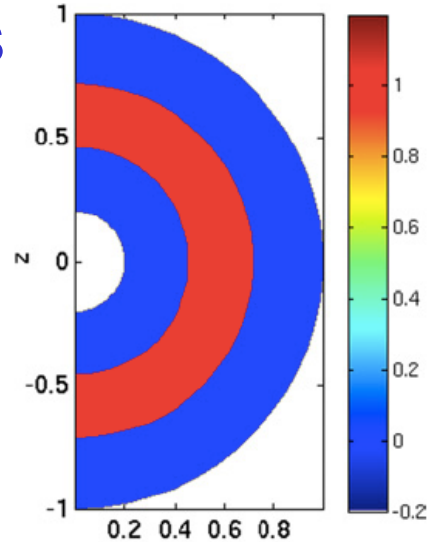
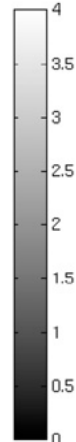
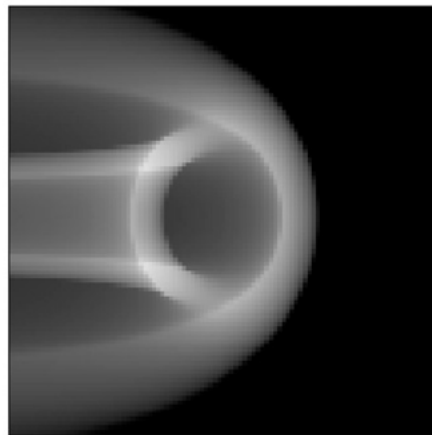


$$I_0(x, y) = \int_{s_C}^{\infty} S_0(\Psi(s), \theta(s), \varphi(s)) ds$$

# Tomography inversion in presence of noise

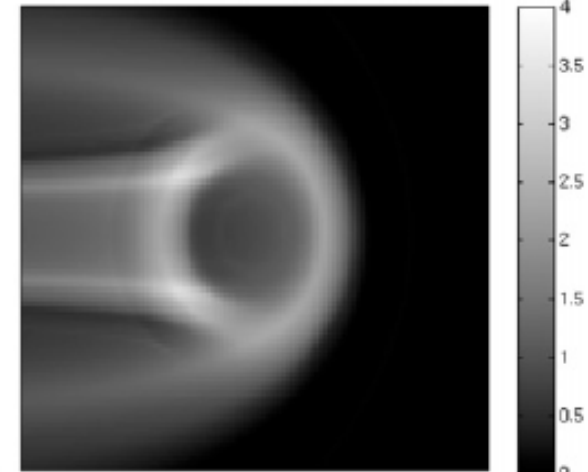
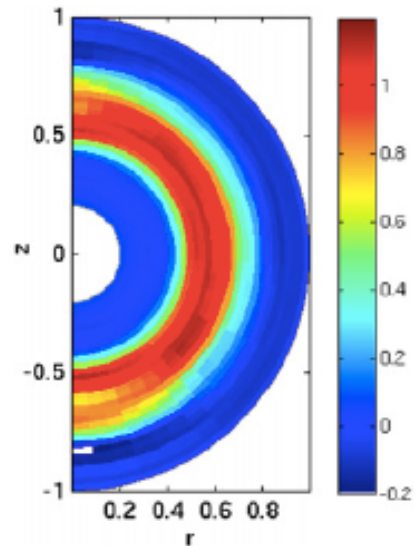
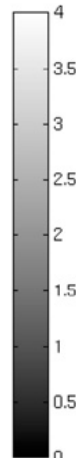
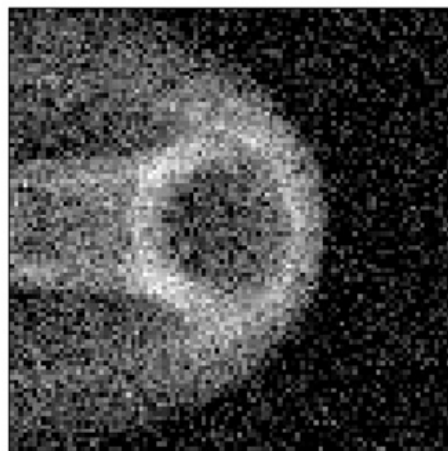
Image received by the camera: integral of the volume emissivity  $I=KS$

Plasma light emissivity  $S$



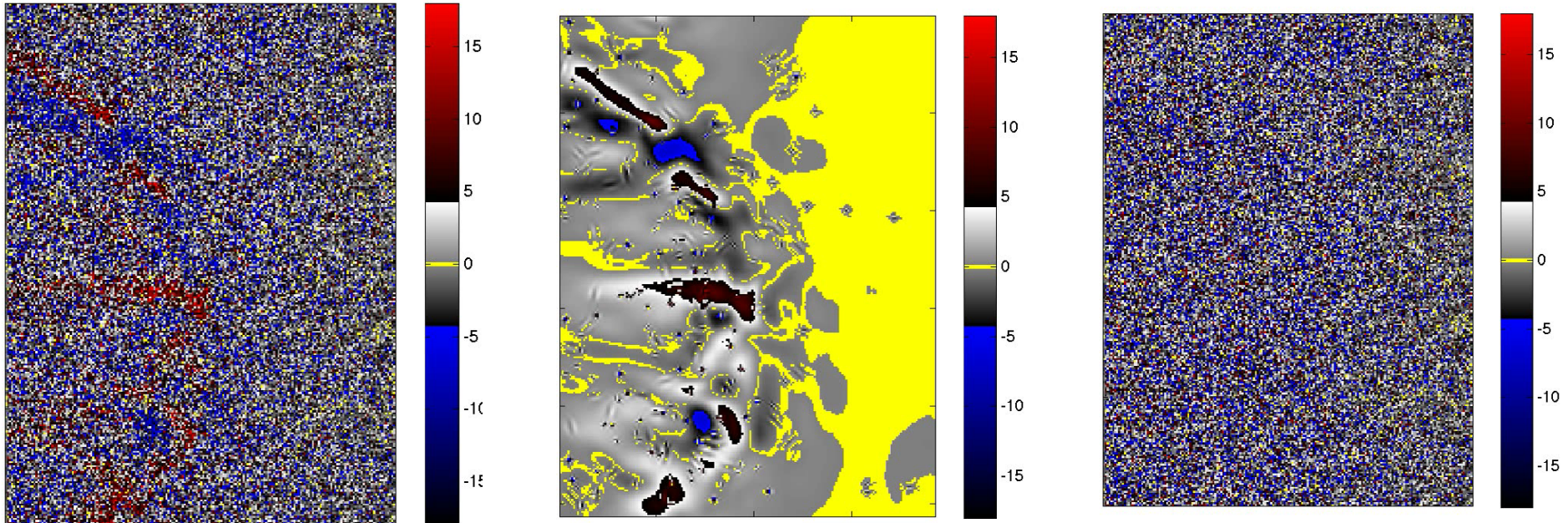
*Nguyen, Fedorczak, Brochard,  
Bonhomme, Schneider, Farge,  
Monier-Garbet, Nuclear Fusion,  
52, 2012*

Denoised plasma emissivity



# Movie from a fast camera in Tore-Supra tokamak

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Noisy  
images

Coherent  
structures

Incoherent  
background

*Nguyen, Fedorczak, Brochard,  
Bonhomme, Schneider, Farge,  
Monier-Garbet, Nuclear Fusion, 52, 2012*



## Noise reduction in plasma simulations using particles

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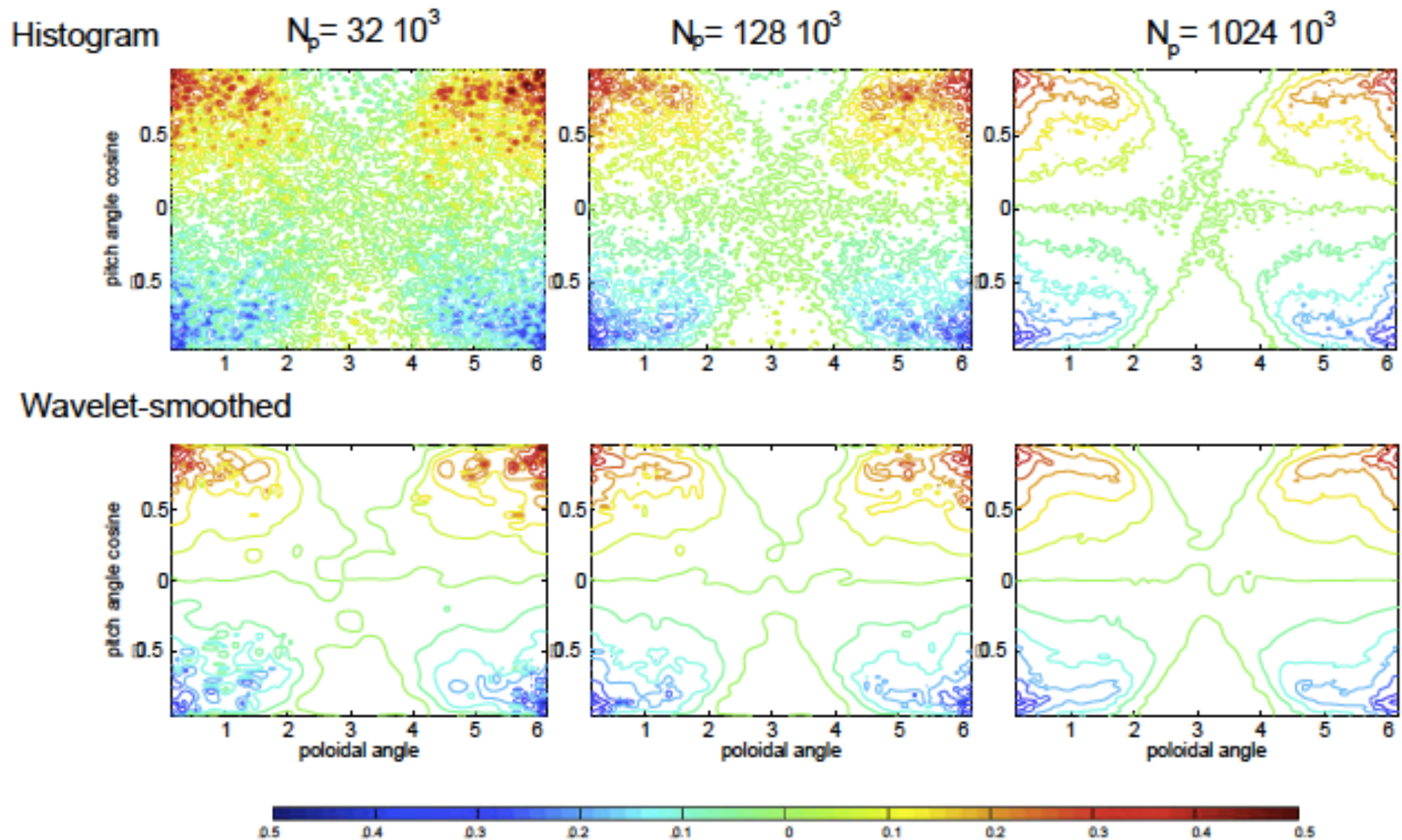
- Accuracy of particle simulations is limited by noise (statistical sampling, not enough particles and grid effects)
- Wavelet based density estimation, accurate estimation of distribution functions with localized sharp features
- Preservation of moments in the distribution functions
- No a priori selection of a global smoothing scale
- No constraints on the dimensionality
- Computationally efficient: same order as for finite size particle approach

*Nguyen van yen, del-Castillo-Negrete,  
Schneider, Farge and Chen,  
J. Comput. Phys., 229, 2010*



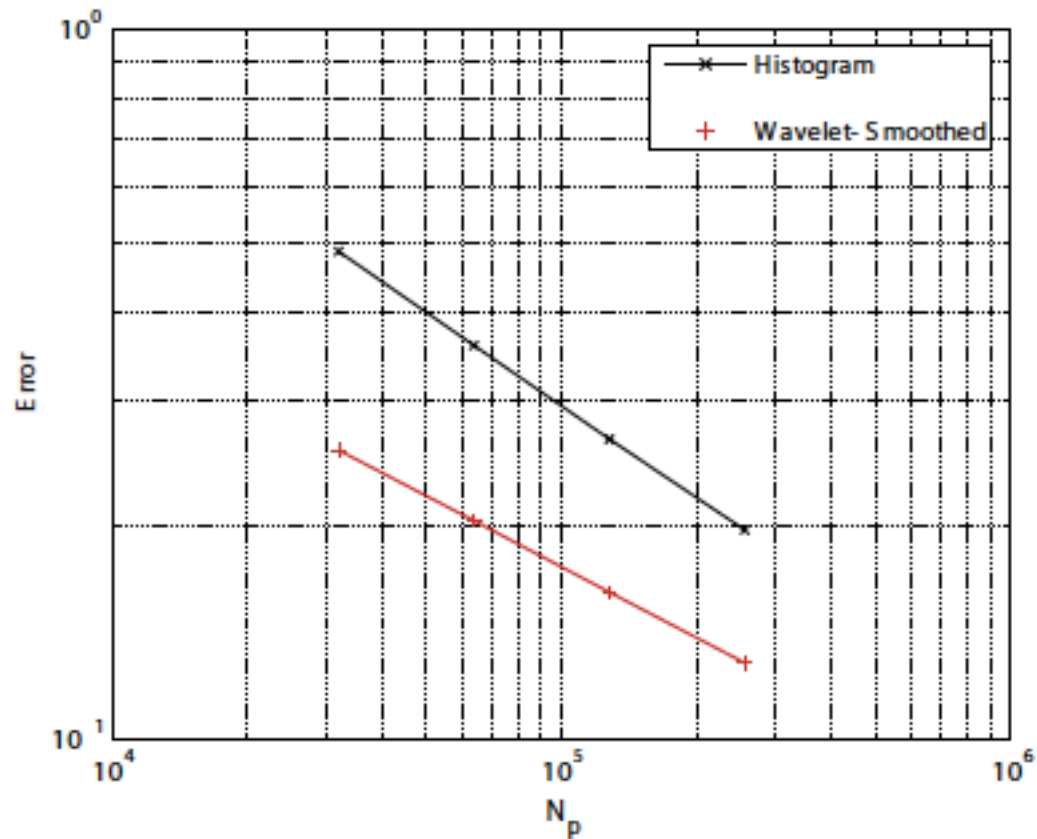
# Noise reduction in plasma simulations using particles

## Collisional guiding center transport data (Delta5d)



# Noise reduction in plasma simulations using particles

## Collisional guiding center transport data (Delta5d)



Nguyen van yen,  
del-Castillo-Negrete,  
Schneider, Farge  
and Chen,  
*J. Comput. Phys.*,  
229, 2010

RMS error estimate with respect to the reference density computed with  $N_p = 1024 \times 10^3$ .

Error reduction by about a factor 2.

## Particle in wavelets scheme for Vlasov-Poisson equation

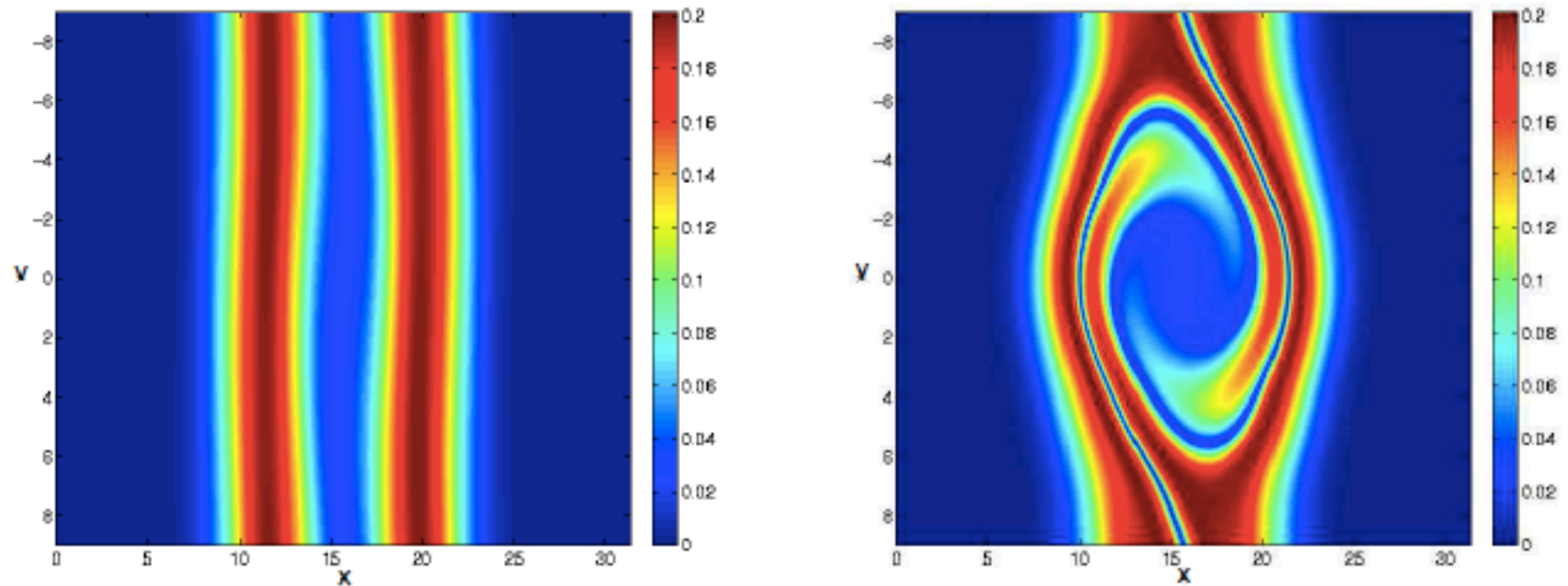
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- Plasma distribution function is discretized using tracer particles
- The charge distribution is reconstructed using wavelet based density estimation
- Wavelet expansion of the Dirac delta functions corresponding to each particle
- Wavelet Galerkin Poisson solver to compute the electric potential from the electron charge density (diagonal preconditioning)
- Improvement of precision compared to a classical PIC scheme for a given number of particles

*Nguyen van yen, Sonnendrücker,  
Schneider and Farge,  
ESAIM Proc., 32, 2011*

## Particle in wavelets scheme for Vlasov-Poisson equation

### Two-stream instability test case

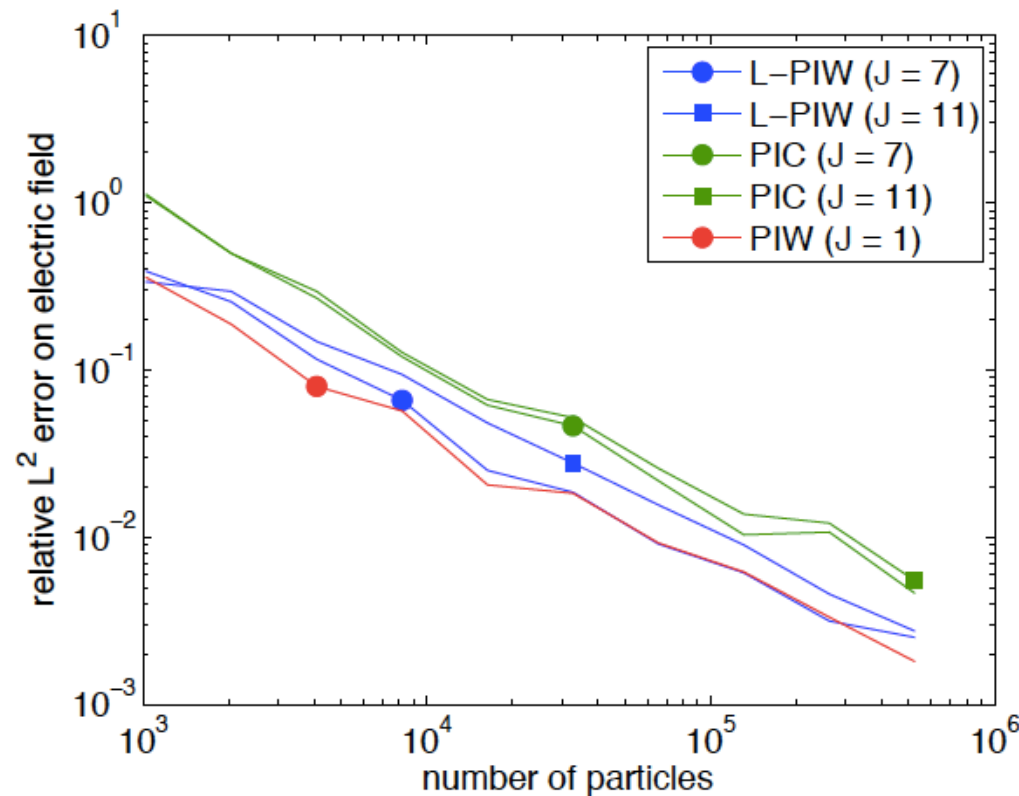


Particle distribution function at  $t=10$  (left) and  $t=30$  (right).

*Nguyen van yen, Sonnendrücker,  
Schneider and Farge,  
ESAIM Proc., 32, 2011*

## Particle in wavelets scheme for Vlasov-Poisson equation

$L^2$  error on the electric field at  $t = 30$ ,  
as a function of number of particles



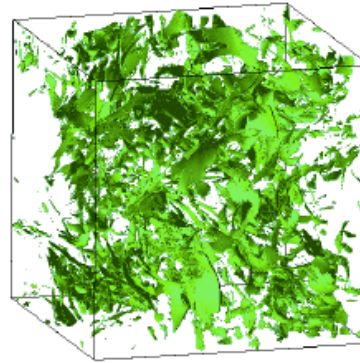
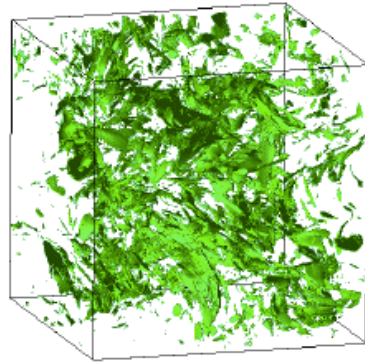
Nguyen van yen, Sonnendrücker,  
Schneider and Farge,  
ESAIM Proc., 32, 2011



# Coherent structures extraction in 3D MHD flow

**Total**

100 % N



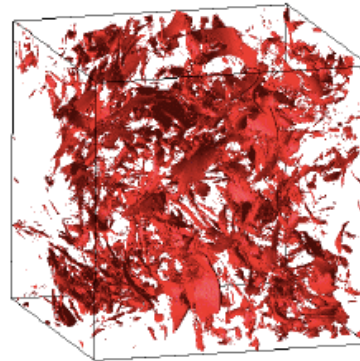
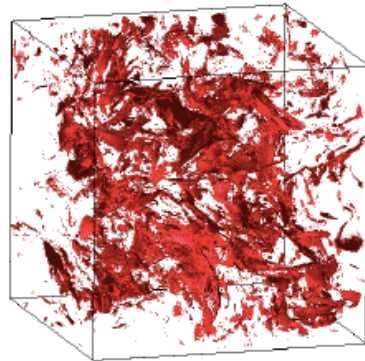
100 % N

**Vorticity**

**Current density**

**Coherent**

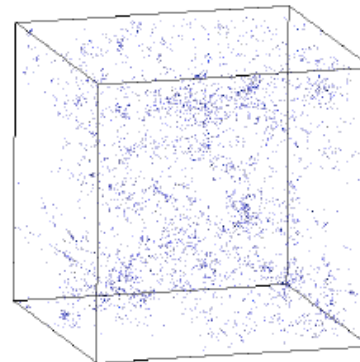
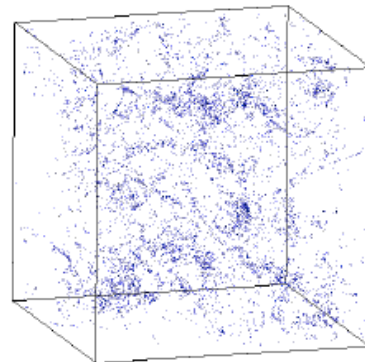
3.21 % N



3.16 % N

**Incoherent**

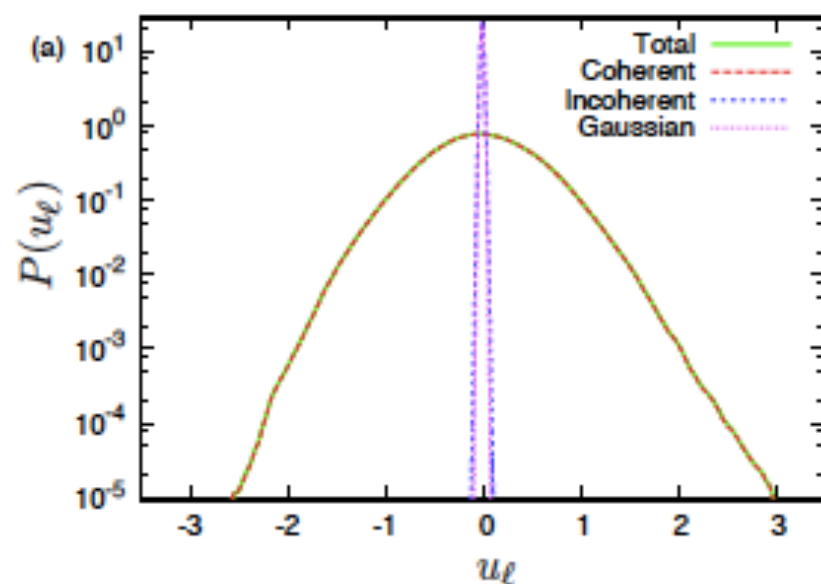
96.79 % N



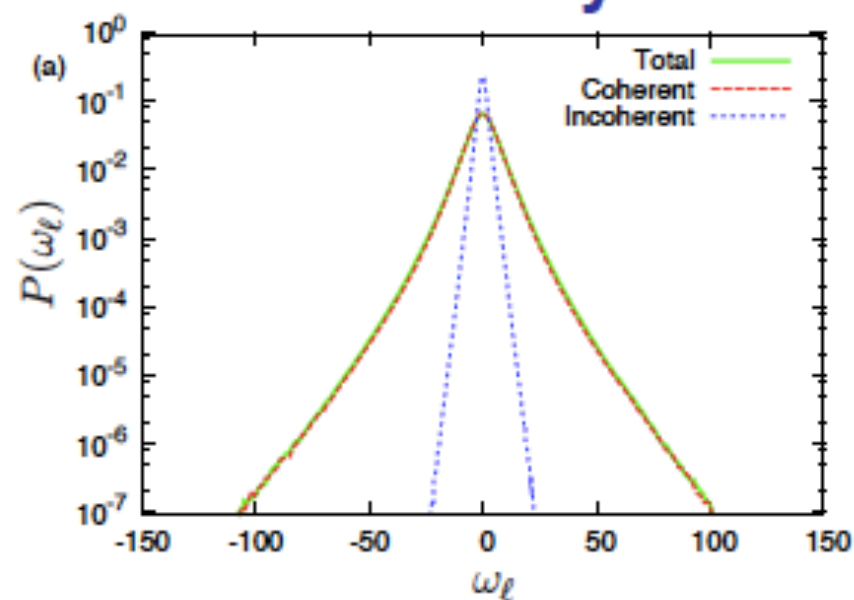
96.84 % N

Yoshimatsu,  
Kondo,  
Schneider,  
Okamoto,  
Hagiwara  
& Farge  
*Phys. Plasmas*,  
16, 2009

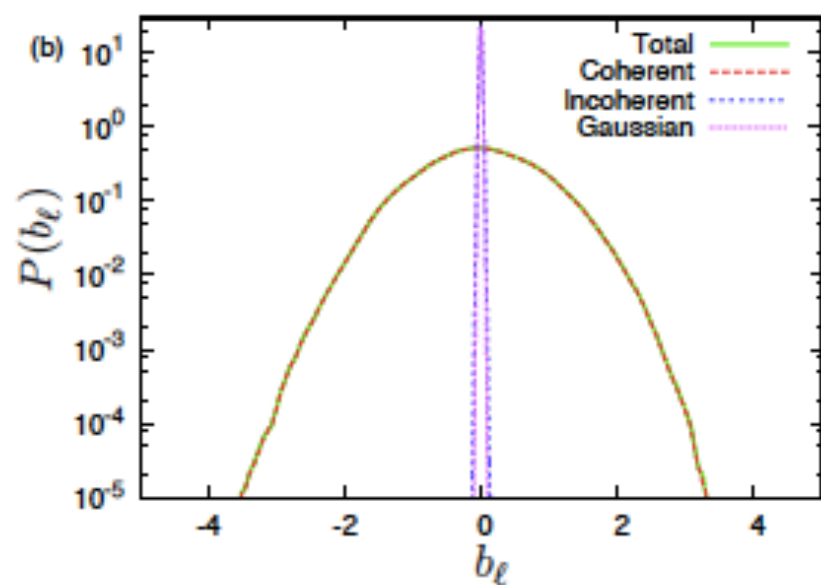
## Velocity



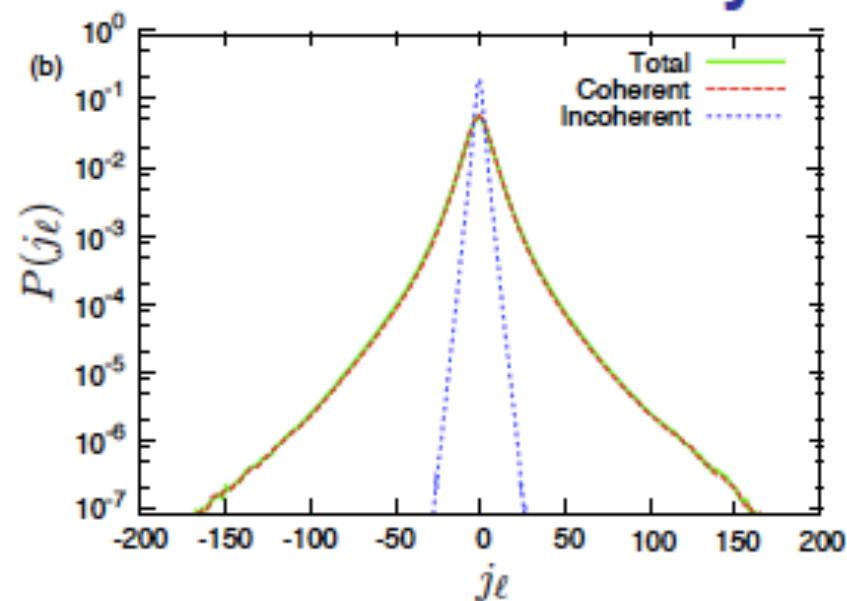
## Vorticity



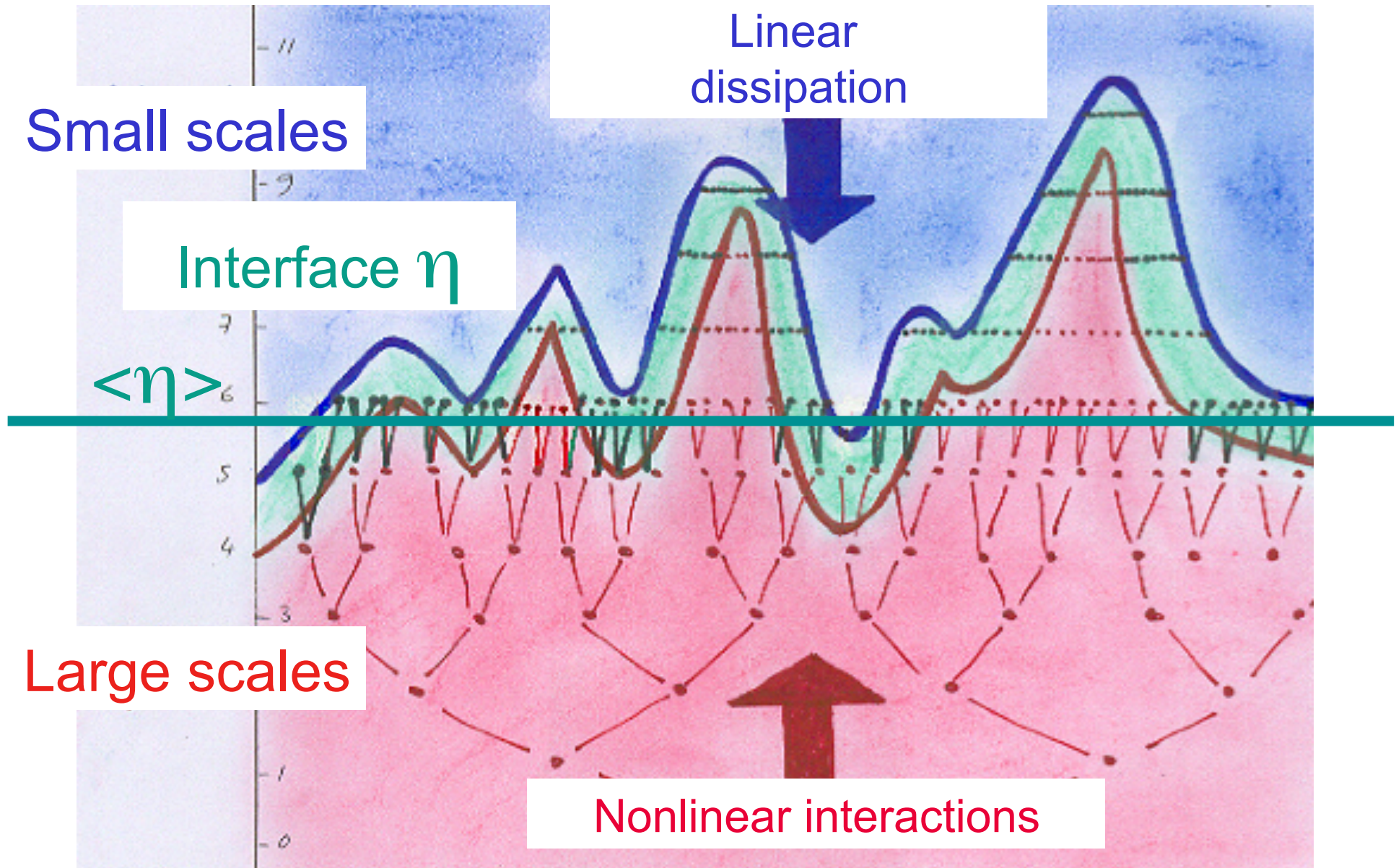
## Magnetic field



## Current density

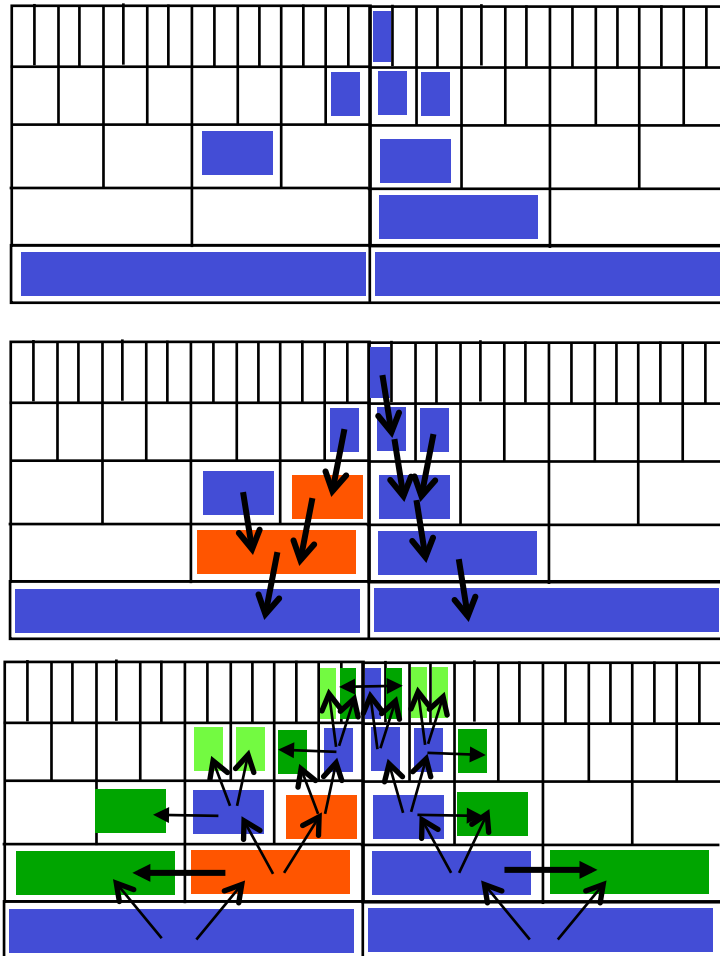


# Coherent Vorticity Simulation (CVS)





# Coherent Vortex Simulation (CVS)



Schneider & Farge, 2000,  
*Comp. Rend. Acad. Sci. Paris*, 328

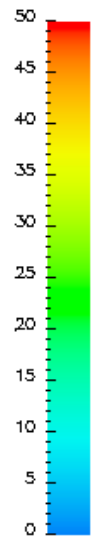
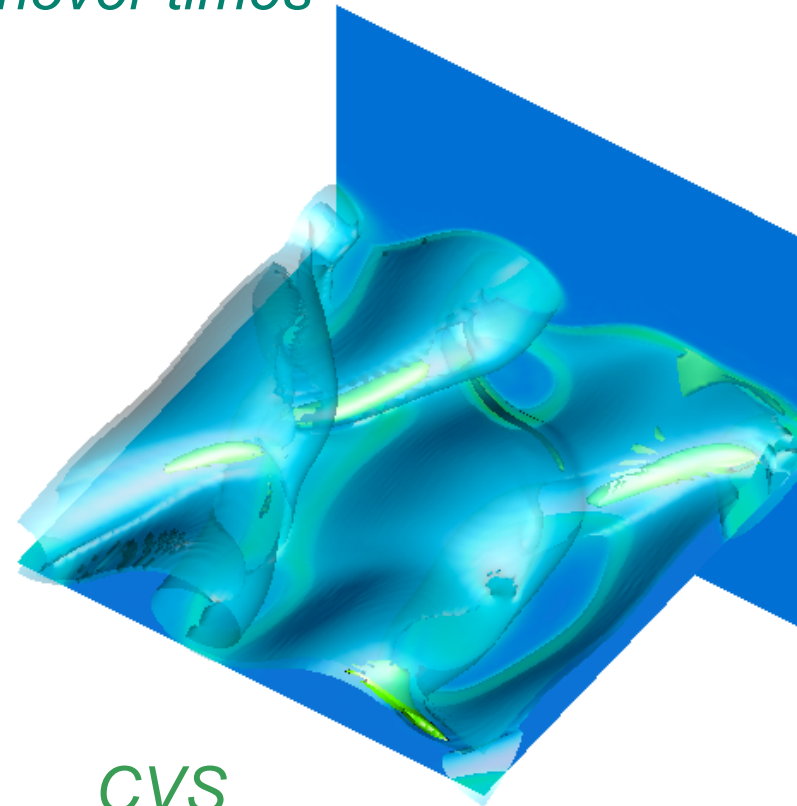
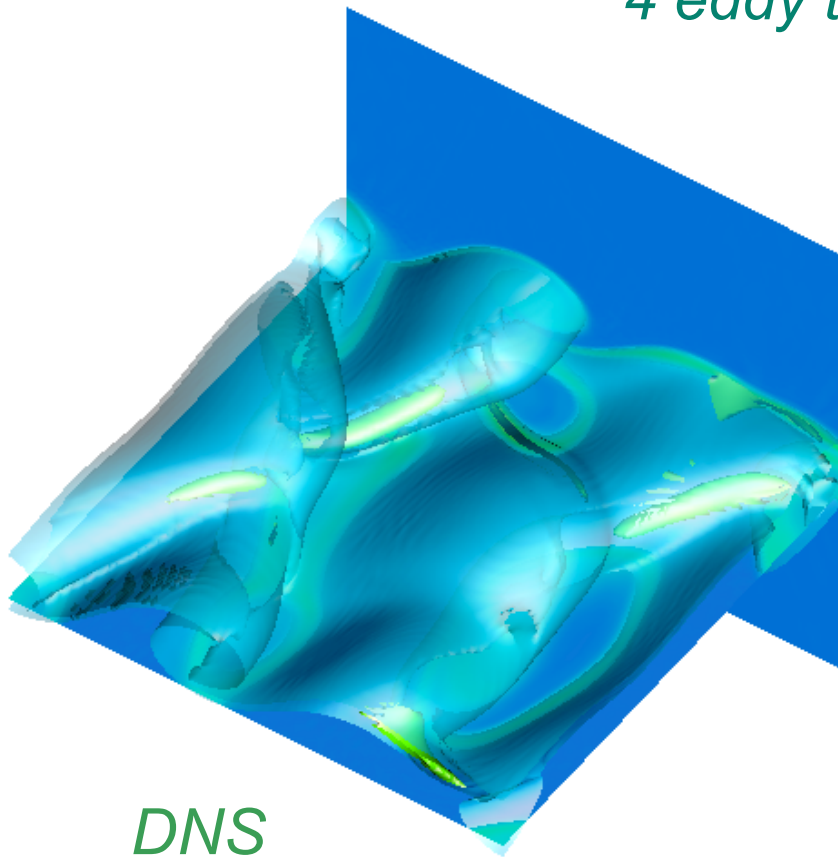
1. Selection of the wavelet coefficients whose modulus is larger than the threshold.
2. Construction of a 'graded-tree' which defines the 'interface' between the coherent and incoherent wavelet coefficients.
3. Addition of a 'security zone' which corresponds to dealiasing.

Schneider & Farge, 2002,  
*Appl. Comput. Harmonic Anal.*, 12

Schneider, Farge et al., 2005,  
*J. Fluid Mech.*, 534(5)

# 3D turbulent mixing layer

4 eddy turnover times

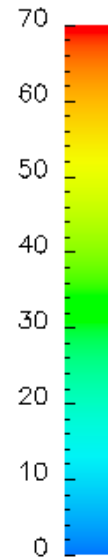
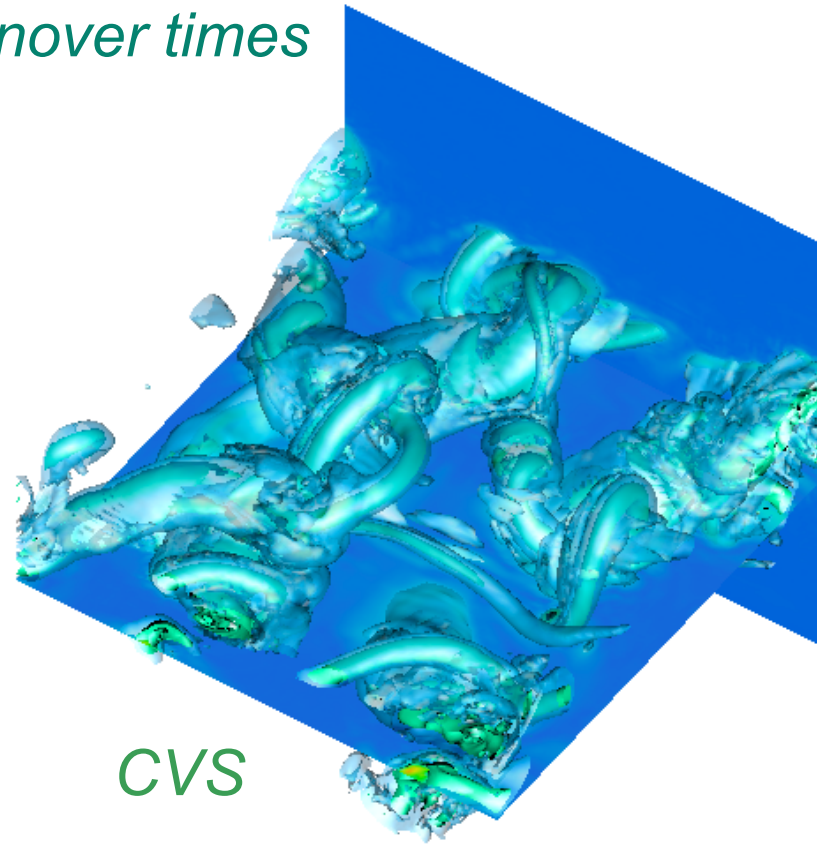
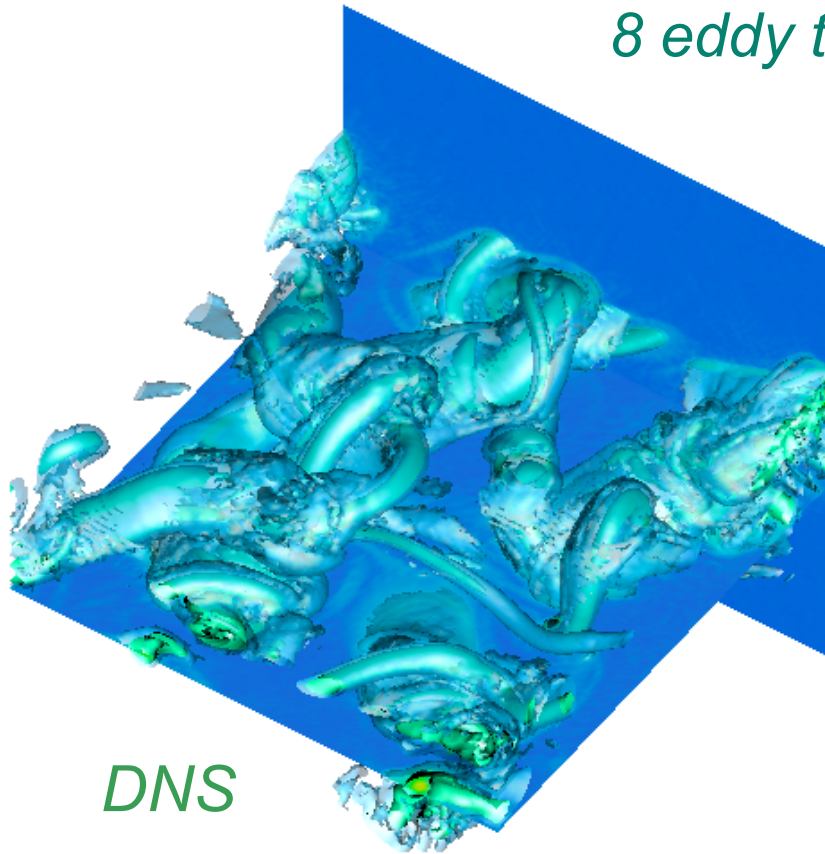


Schneider, Farge,  
Pellegrino, Rogers 2005,  
*J. Fluid Mech.*, 534(5)



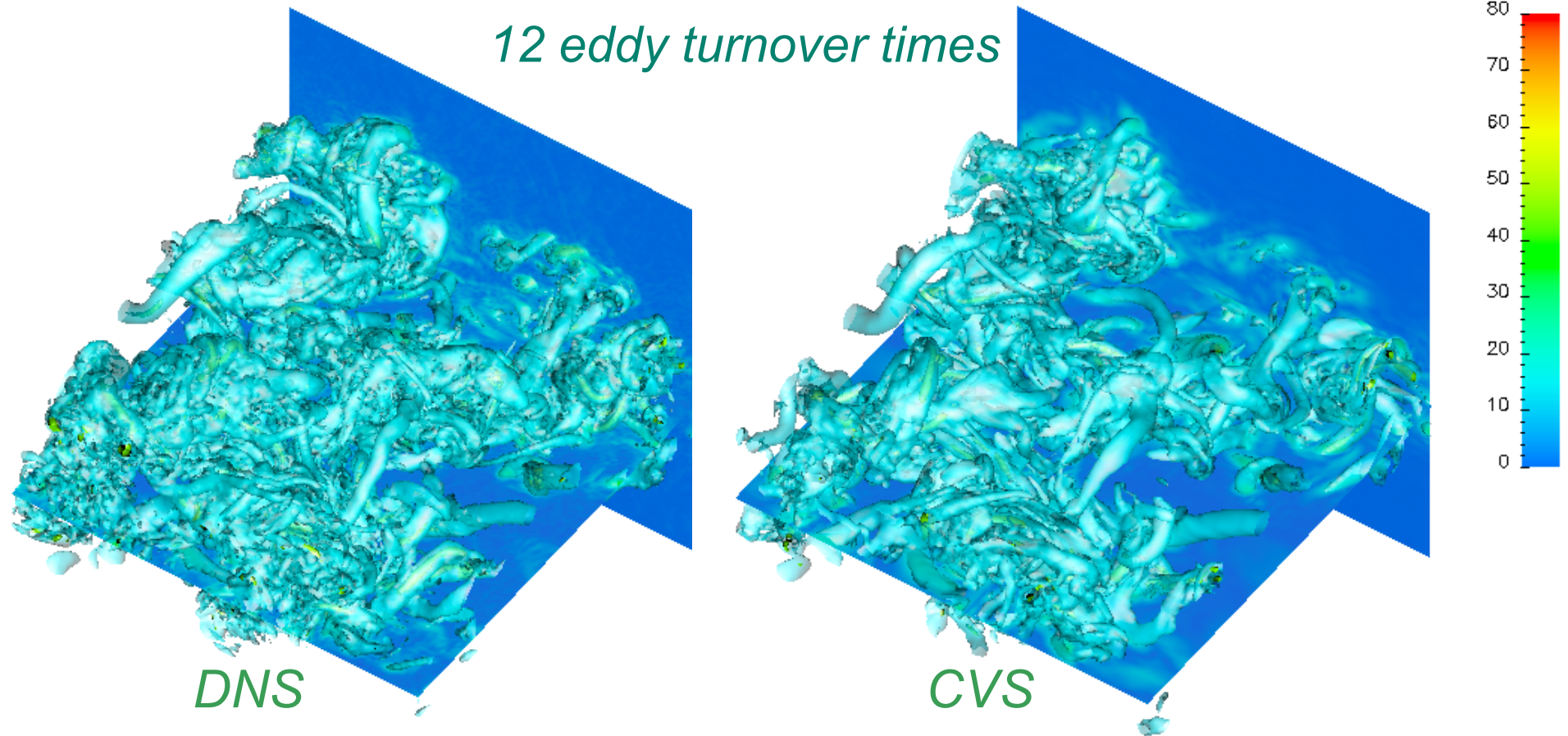
# 3D turbulent mixing layer

*8 eddy turnover times*



Schneider, Farge,  
Pellegrino, Rogers 2005,  
*J. Fluid Mech.*, 534(5)

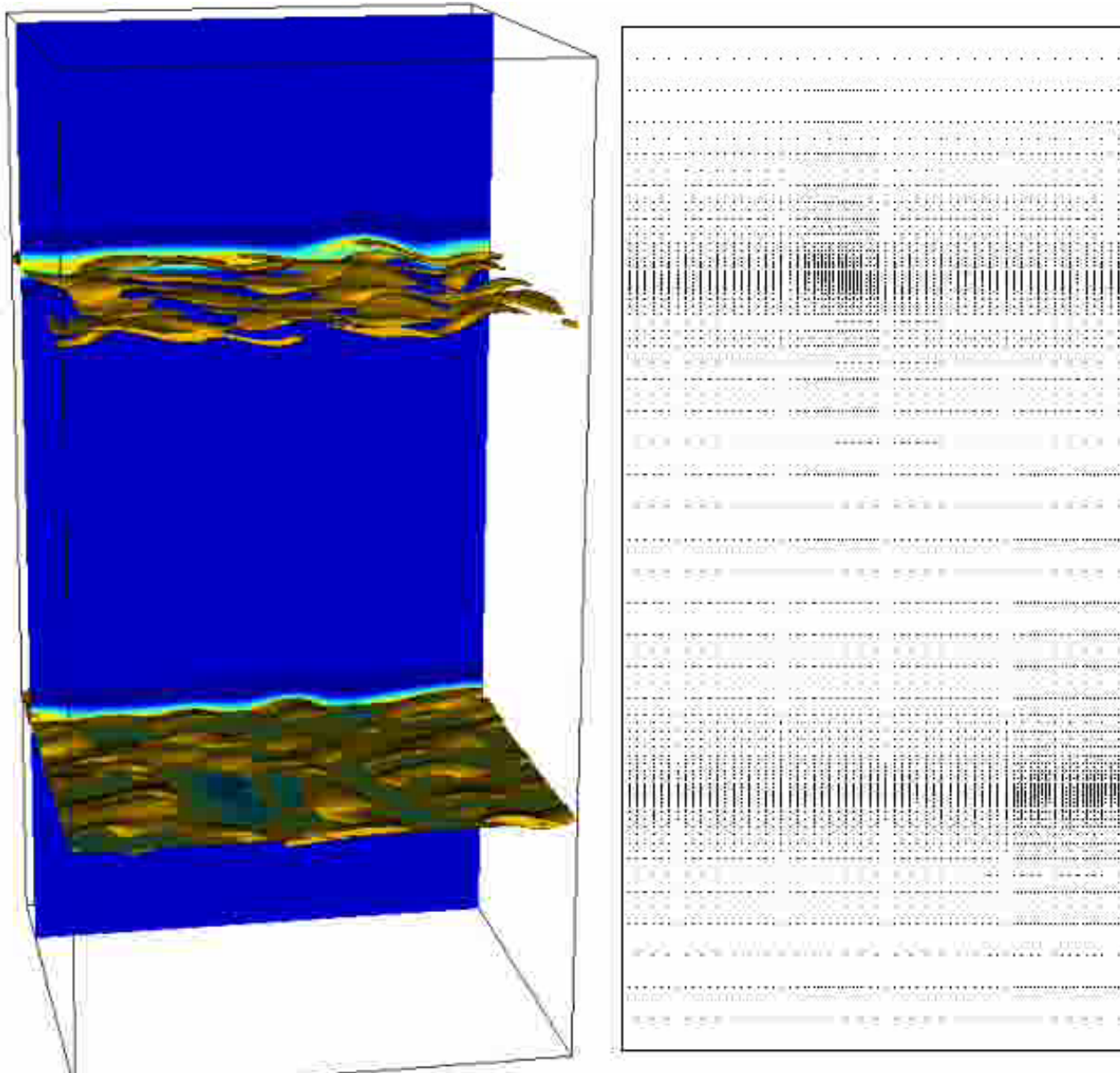
# 3D turbulent mixing layer



Schneider, Farge,  
Pellegrino, Rogers 2005,  
*J. Fluid Mech.*, 534(5)

# Adaptive computation using wavelets

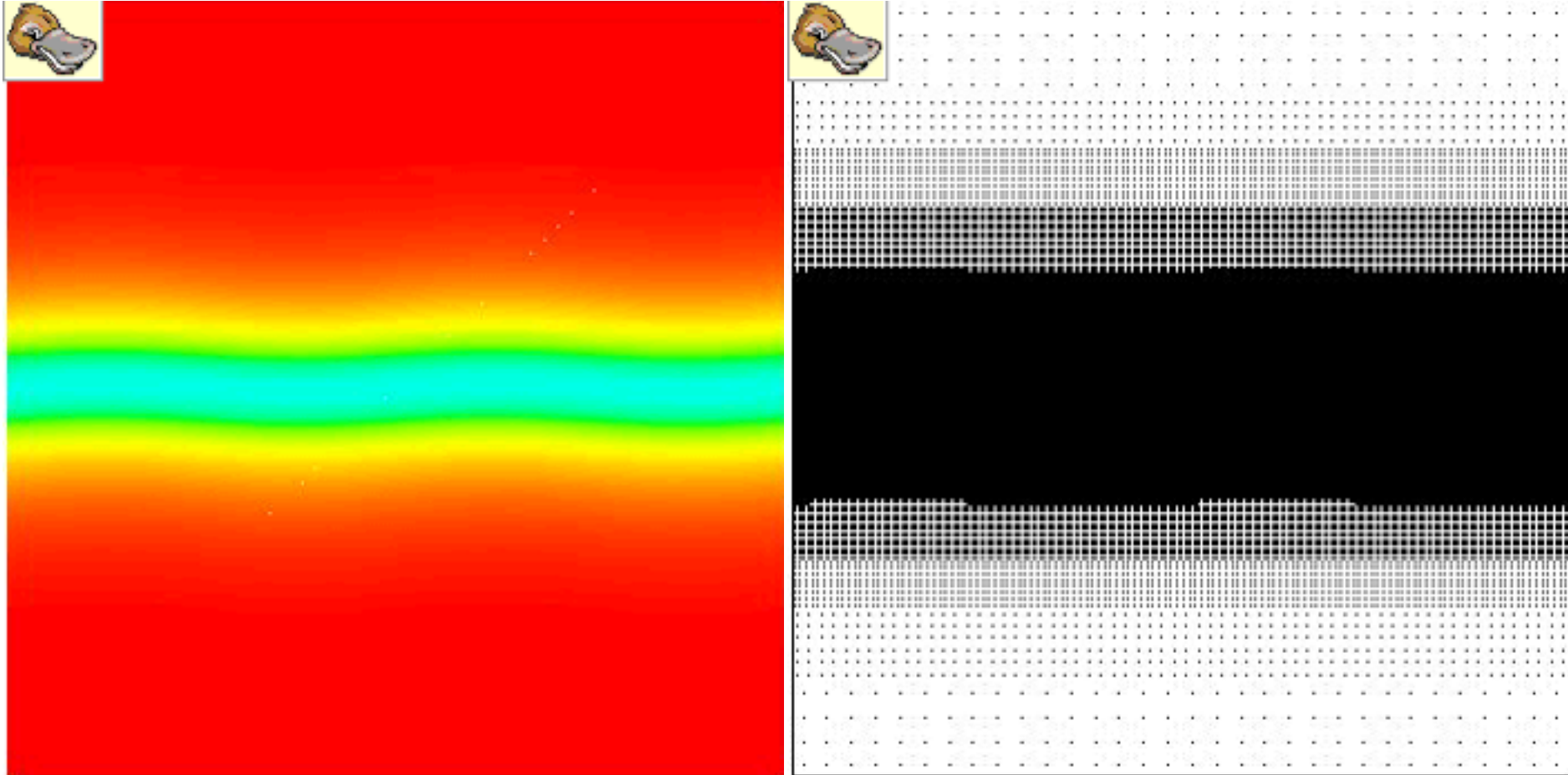
*Koster, Schneider, Griebel, Farge  
Numerical Flow Simulation II,  
75, Springer, 2001*





# Adaptive computation using wavelets

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*Roussel and Schneider,  
75, 2000*

# Turbulence practice is the 'art of averaging'

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Reynolds averaging (1883) :

$$\text{Field } f = \text{Mean } \bar{f} + \text{Fluctuations } f'$$

$$\text{with } \overline{f'} = 0 \quad \overline{\bar{f}} = \bar{f}$$

$$\overline{f + g} = \bar{f} + \bar{g} \quad \overline{\partial f} = \partial \bar{f}$$

but nonlinearity is hard to handle since there is no scale separation :

$$\overline{fg} = \bar{f}\bar{g} + \overline{f'g'}$$

New way of averaging (1992):

$$f' = f'_c + f'_i$$

*Fluctuations* = *coherent fluctuations* + *incoherent fluctuations*  
= *intermittent fluctuations* + *non-intermittent fluctuations*



# Review papers on wavelets

## <http://wavelets.ens.fr>

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Marie Farge, 1992

Wavelet Transforms and Their Applications to Turbulence

*Ann. Rev. Fluid Mech.*, **24**, 395-457

Marie Farge, Nicholas Kevlahan, Valerie Perrier and Eric Goirand, 1996

Wavelets and Turbulence

*IEEE Proceedings*, **84**, 4, 1996, 639-669

Marie Farge, Nicholas Kevlahan, Valérie Perrier and Kai Schneider, 1999

Turbulence Analysis, Modelling and Computing using Wavelets

*Wavelets in Physics*, ed. J. van den Berg, Cambridge University Press, 117-200

Marie Farge and Kai Schneider, 2002

Analyzing and computing turbulent flows using wavelets

*Summer Course, Les Houches LXXIV, New trends in turbulence*, Springer

Kai Schneider and Marie Farge, 2006

Wavelets: Mathematical Theory

*Encyclopedia of Mathematical Physics*, Elsevier, 426-437

Marie Farge and Kai Schneider, 2006

Wavelets: Application to Turbulence

*Encyclopedia of Mathematical Physics*, Elsevier, 408-419

# Papers on applications to tokamaks

## <http://wavelets.ens.fr>

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Marie Farge, Kai Schneider and Pascal Devynck, 2006

Extraction of coherent bursts in turbulent edge plasma using orthogonal wavelets

*Physics of Plasmas*, **13**(2), 042304, 1-11

Romain Nguyen van yen, Diego del Castillo–Negrete, Kai Schneider, Marie Farge and Guangye Chen, 2010

Wavelet–based density estimation for noise reduction in plasma simulation using particles

*J. Comput. Phys.*, **229**(8), 2821-2839

Romain Nguyen van yen, Eric Sonnendrücker, Kai Schneider and Marie Farge, 2011

Particle-in-wavelet scheme for the 1D Vlasov-Poisson equations

*ESAIM Proc.*, **32**, 134-148

Romain Nguyen van yen, Nicolas Fedorczak, Frédéric Brochard, Kai Schneider, Marie Farge and Pascale Monier-Garbet, 2012

Tomographic reconstruction of tokamak edge turbulence light emission from single image using wavelet-vaguelette decomposition

*Nuclear Fusion, IAEA (International Atomic Energy Agency)*, **52**, 013005, 1-11

<http://wavelets.ens.fr>

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